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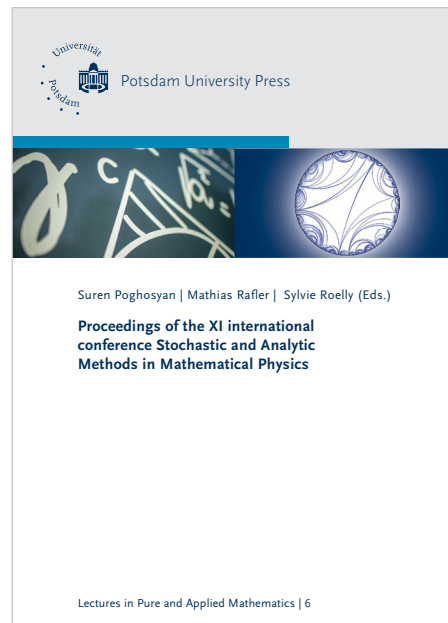
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Construction of limiting Gibbs processes and the uniqueness of Gibbs processes

Suren Poghosyan* and Hans Zessin†

Abstract. For a pair potential Φ in a general phase space X satisfying some natural and sufficiently general stability and regularity conditions in the sense of Poghosyan and Ueltschi we define by means of the so-called Ursell kernel a function r which is shown to be the correlation function of a unique infinitely extended process P . Finally, under more restrictive assumptions, we show that the Gibbs process for Φ , if it exists, coincides with P . Here we use the classical method of Kirkwood-Salsburg equations.

1 Preliminaries

Let $(X, \mathcal{B}(X), \mathcal{B}_0(X))$ be the underlying phase space where X is a locally compact, second countable Hausdorff topological space, $\mathcal{B}(X)$ its Borel σ -field and $\mathcal{B}_0(X)$ its bounded Borel sets. Let ρ be a Radon measure on X .

Let $\mathcal{M}^{\cdot}(X)$ be the space of Radon point measures on X and \mathfrak{X} be the collection of all finite point measures (finite configurations) ξ in X . \mathfrak{X}_+ denotes the collection of all

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non-empty ξ in \mathfrak{X} . For $A \in \mathcal{B}_0(X)$ let $\mathfrak{X}(A)$ be the set of finite point measures supported by A . Let $\mathcal{M}_R(X)$, $R > 0$, be the space of simple point measures μ on X having the property that the minimal distance of every pair of points in any configuration μ is R , i. e. $(x, y \in \mu, x \neq y \Rightarrow d(x, y) > R)$ where d is a metric on X . Here all Dirac measures ε_x at the point $x \in X$ and the zero measure o are elements of \mathfrak{X} and $\mathcal{M}_R(X)$.

We call a subset \mathfrak{X}' of \mathfrak{X} an *environment in X* if $(\eta \in \mathfrak{X}', \xi \preceq \eta \Rightarrow \xi \in \mathfrak{X}')$. Here $\xi \preceq \eta$ if $\xi(x) \leq \eta(x)$ for all $x \in X$. Examples are \mathfrak{X} and $\mathfrak{X}_R = \mathcal{M}_R(X) \cap \mathfrak{X}$.

We denote by F_+ the space of $[0, +\infty]$ -valued measurable functions on the corresponding space and by \mathcal{K} we denote the collection of continuous functions with compact support. Define a locally finite measure Λ_ρ on \mathfrak{X} by

$$\Lambda_\rho \varphi = \varphi(o) + \sum_{n=1}^{\infty} \frac{1}{n!} \int_X \cdots \int_X \varphi(\varepsilon_{x_1} + \cdots + \varepsilon_{x_n}) \rho(dx_1) \cdots \rho(dx_n), \quad \varphi \in F_+.$$

For a given configuration $\mu \in \mathcal{M}^{\cdot\cdot}$ we define the following measure on \mathfrak{X} :

$$\begin{aligned} \Lambda'_\mu(h) &= \sum_{n=0}^{\infty} \frac{1}{n!} \int_{X^n} h(\varepsilon_{x_1} + \cdots + \varepsilon_{x_n}) \tilde{\mu}^n(dx_1, \dots, dx_n), \quad h \in F_+, \text{ where} \\ \tilde{\mu}^n(dx_1, \dots, dx_n) &= (\mu - \varepsilon_{x_1} - \cdots - \varepsilon_{x_{n-1}})(dx_n) \cdots (\mu - \varepsilon_{x_1})(dx_2) \mu(dx_1). \end{aligned}$$

$\tilde{\mu}^n$ is called the *factorial measure* of μ of order n , and Λ'_μ the *compound factorial measure built on μ* . The term $n = 0$ of the sum is $h(o)$. Also, $\Lambda'_o(h) = h(o)$.

Below we often use the following important equation, the *Minlos' formula* [3]:

$$\int_{\mathfrak{X}} \int_{\mathfrak{X}} h(\xi, \nu - \xi) \Lambda'_\nu(d\xi) \Lambda_\rho(d\nu) = \int_{\mathfrak{X}} \int_{\mathfrak{X}} h(\xi, \nu) \Lambda_\rho(d\xi) \Lambda_\rho(d\nu), \quad h \in F_+, \quad (5.1)$$

which is valid for all h integrable with respect to the measure on the left-hand or the right-hand side of the equation.

Let P be a point process in X that is a probability on $\mathcal{M}^{\cdot\cdot}(X)$. The *moment measure of P of order k* is the measure on X^k defined by

$$\nu_P^k f = \int_{\mathcal{M}^{\cdot\cdot}(X)} \mu^{\otimes k}(f) P(d\mu), \quad f \in \mathcal{K}(X^k),$$

whereas the *correlation measure* (also called *factorial moment measure*) of P of order k is the measure given by

$$\tilde{\nu}_P^k(f) = \int_{\mathcal{M}^{\cdot}(X)} \tilde{\mu}^k(f) P(d\mu), \quad f \in \mathcal{H}(X^k).$$

If $\tilde{\nu}_P^k$ has a density r_P^k with respect to some product measure $\rho^{\otimes k}$, where ρ is a Radon measure on X , then we say that r_P^k is a *correlation function of P of k -th order*. The process P is called *of order k* if ν_P^k is a Radon measure. P is called *of infinite order* if it is of order k for every k .

2 Ruelle's algebraic approach

We here follow Ruelle [10]. Let \mathcal{A} be the set of all measurable complex functions on \mathfrak{X} . We define a \star -multiplication of two functions $h_1, h_2 \in \mathcal{A}$ by

$$h_1 \star h_2(\xi) = \int_{\mathfrak{X}} h_1(v) h_2(\xi - v) \Lambda'_\xi(dv), \quad \xi \in \mathfrak{X}. \quad (5.2)$$

With the \star -product \mathcal{A} becomes a commutative algebra with the unit $\mathbb{1}(\xi) = \delta_o(\xi)$. Let $\mathcal{A}_0 = \{f \in \mathcal{A} \mid f(o) = 0\}$. We define the mapping $\Gamma : \mathcal{A}_0 \rightarrow \mathbb{1} + \mathcal{A}_0$ (*algebraic exponent*) by

$$\Gamma h = \mathbb{1} + h + \frac{1}{2!} h^{\star 2} + \dots + \frac{1}{n!} h^{\star n} + \dots, \quad h \in \mathcal{A}_0. \quad (5.3)$$

Let Φ be a measurable symmetric function $\Phi : X \times X \rightarrow]-\infty, +\infty]$, a pair potential in X . $E(\xi) := \sum_{1 \leq i < j \leq n} \Phi(x_i, x_j)$ is the *energy* of the configuration $\xi = \varepsilon_{x_1} + \dots + \varepsilon_{x_n}$; $B := e^{-E}$ is called the Boltzmann factor. The *conditional energy* at x given the configuration ξ is given by $W_\Phi(x, \xi) := \int_X \Phi(x, y) \xi_x(dy)$, where $\xi_x = \xi$ if $x \notin \xi$ and $\xi_x = \xi - \varepsilon_x$ otherwise.

We assume that Φ is *b-stable*, i. e. there exists a measurable function $b : X \rightarrow [0, +\infty)$ such that $E(\xi) \geq -\sum_{x \in \xi} b(x)$, $\xi \in \mathfrak{X}$.

We consider also \mathcal{P} -stable¹ Φ with *stability function* b in the environment \mathfrak{X}' . This means that there exists a measurable function $b : X \rightarrow [0, +\infty)$ such that $W_\Phi(x, \xi_x) \geq -b(x)$, $x \in \xi \in \mathfrak{X}'$, $\xi_x = \xi - \varepsilon_x$. If Φ is \mathcal{P} -stable with function b , then it is *b-stable*.

Any non-negative Φ is \mathcal{P} -stable in the environment \mathfrak{X} . Another important example is the *Penrose potential* [7] (see also [5]). Let (X, ρ) be the d -dimensional Euclidean

¹This notion goes back to Oliver Penrose [7].

space with Lebesgue measure. Let $c, \varepsilon, R > 0$ be constants. Φ is the following hard-core potential: If $|x - y| < R$ then $\Phi(x, y) = +\infty$; and if $|x - y| \geq R$ then $|\Phi|(x, y) \leq c|x - y|^{-(d+\varepsilon)}$. As Penrose has shown this potential is \mathcal{P} -stable in the environment \mathfrak{X}_R with a constant stability which can be calculated explicitly.

The Boltzmann factor $B = e^{-E}$ is an element of the algebra (\mathcal{A}, \star) having an inverse with respect to the \star multiplication, which is denoted by B_\star^{-1} . Another important element U of the algebra \mathcal{A} is the *Ursell function* given by

$$U(o) = 0, \quad U(\varepsilon_x) = 1, \quad U(\varepsilon_{x_1} + \dots + \varepsilon_{x_n}) = \sum_{\gamma \in \mathcal{C}_n} \prod_{(i,j) \in \gamma} \omega(x_i, x_j), \quad n \geq 2, \quad (5.4)$$

where \mathcal{C}_n denotes the set of all simple, unoriented, connected graphs γ with n vertices, the product is taken over all edges (i, j) in γ and $\omega(x, y) = e^{-\Phi(x,y)} - 1$ is the Mayer function.

Note that $U \in \mathcal{A}_0$ and the following important relation is valid $B = \Gamma U$.

3 Ursell kernel Representation of the correlation function

Here we follow the work of Minlos, Poghosyan [4]. Let $z : X \rightarrow [0, +\infty)$ be measurable. We consider Radon measures of the form $z \cdot \rho = \rho_z$, where ρ is Radon measure and z is a density function.

Given $A \in \mathcal{B}_0(X)$ we define the finite volume Gibbs process in A as the probability $Q_{z,A}$ on $\mathfrak{X}(A)$ which is given by

$$Q_{z,A}(d\xi) = \frac{1}{\Xi(z,A)} e^{-E(\xi)} \cdot \Lambda_{z \cdot \rho_A}(d\xi)$$

where $\rho_A = 1_A \cdot \rho$ and the normalising constant (*the partition function*) is given by

$$\Xi(z,A) = \int_{\mathfrak{X}(A)} \prod_{x \in \eta} z(x) e^{-E(\eta)} \Lambda_{\rho_A}(d\eta).$$

By stability $\Xi(z,A) \leq \exp\left(\int_A e^{b(x)} z(x) \rho(dx)\right) < \infty$.

It is well known that the *correlation function of the Gibbs process* $Q_{z,A}$ is given by

$$r_{z,A}(\xi) = \frac{\prod_{x \in \xi} z(x)}{\Xi(z,A)} \int_{\mathfrak{X}(A)} e^{-E(\xi+\eta)} \Lambda_{z,\rho_A}(d\eta), \quad \xi \in \mathfrak{X}(A).$$

Proposition 5.1 The correlation function has the following remarkable representation:

$$r_{z,A}(\xi) = \prod_{x \in \xi} z(x) \int_{\mathfrak{X}(A)} G(\xi, \eta) \Lambda_{z,\rho_A}(d\eta), \quad \xi \in \mathfrak{X}(A) \quad (5.5)$$

where the *Ursell kernel* $G : \mathfrak{X}^2 \rightarrow \mathbb{R}$ is given by $G(\xi, \eta) = (B_\star^{-1} \star D_\xi B)(\eta)$, $\xi, \eta \in \mathfrak{X}$ and $D_\xi B(v) = B(\xi + v)$, $v \in \mathfrak{X}$. In particular $G(\varepsilon_x, \eta) = U(\varepsilon_x + \eta)$ where U is the Ursell function.

For the proof we note that by the Minlos' formula

$$\frac{1}{\Xi(z,A)} \Lambda_{z,\rho_A}(D_\xi B) = \frac{1}{\Xi(z,A)} \Lambda_{z,\rho_A}(B \star B_\star^{-1} \star D_\xi B) = \Lambda_{z,\rho_A}(B_\star^{-1} \star D_\xi B).$$

For a given pair potential Φ let $\bar{\Phi} = \Phi$ if Φ is finite and $\bar{\Phi} = 1$ if $\Phi = +\infty$. Let a, b, c be non-negative functions on X . We will say that Φ satisfies

◇ *c-regularity*, if there exists a function a such that

$$\int_X |\omega|(x,y) e^{(c+a)(y)} \rho_z(dy) \leq a(x), \quad x \in X. \quad (5.6)$$

◇ *Modified b-regularity*, if there exists a function a such that

$$\int_X |\bar{\Phi}|(x,y) e^{b(y)+a(y)} \rho_z(dy) \leq a(x), \quad x \in X. \quad (5.7)$$

Both assumptions (5.6) and (5.7) are introduced in [8].

Theorem 5.2 Let Φ be a b -stable pair interaction. Assume also that Φ is $2b$ -regular for a . Then the function

$$r_z(\xi) = \prod_{x \in \xi} z(x) \int_{\mathfrak{X}} G(\xi, \eta) \Lambda_{\rho_z}(d\eta), \quad \xi \in \mathfrak{X} \quad (5.8)$$

is well defined and satisfies the following *Ruelle bound*

$$r_z(\xi) \leq \prod_{x \in \xi} z(x) \int_{\mathfrak{X}} |G|(\xi, \eta) \Lambda_{\rho_z}(d\eta) \leq \prod_{x \in \xi} z(x) e^{(2b+a)(x)}, \quad \xi \in \mathfrak{X}. \quad (5.9)$$

If Φ is \mathcal{P} -stable and b -regular for a then (5.9) holds with $e^{(b+a)(x)}$ instead of $e^{(2b+a)(x)}$. Moreover $r_z(\xi) = \lim_{A \uparrow X} r_{z,A}(\xi)$.

The proof of this theorem is based on the so-called *forest graph estimate*. For $\xi, \eta \in \mathfrak{X}$ let $\mathcal{F}(\xi, \eta)$ be the collection of forests with the set of vertices $\xi + \eta$ and roots ξ . An unoriented simple graph is called *rooted forest* if its connected components are *rooted trees*, i. e. trees where one vertex is specified as a root.

We consider the case of b -stable Φ . The \mathcal{P} -stable case is entirely the same, one only needs to replace e^{2b} by e^b . If Φ is modified regular, then one has to pass from ω to $\bar{\Phi}$ using the formula $|\omega|(x, y) \leq |\bar{\Phi}|(x, y) e^{\Phi^-(x, y)}$.

Lemma 5.3 ([4]) For $\xi \neq \emptyset$,

$$|G|(\xi, \eta) \leq \prod_{x \in \xi + \eta} e^{2b(x)} \sum_{\gamma \in \mathcal{F}(\xi, \eta)} \prod_{(x, y) \in \gamma} |\omega|(x, y). \quad (5.10)$$

Denoting the right-hand side of (5.10) by $H(\xi, \eta)$ one can show that

$$H(\varepsilon_{x_1} + \dots + \varepsilon_{x_n}, \eta) = H(x_1, \cdot) \star \dots \star H(x_n, \cdot)(\eta).$$

Then an application of the Minlos' formula and Theorem 2.1 from [8] completes the proof of Theorem 5.2.

In particular Lemma 5.3 gives the famous tree graph estimate of the Ursell function:

$$|U|(\eta) = |G|(\varepsilon_x, \eta - \varepsilon_x) \leq \prod_{x \in \eta} e^{2b(x)} \sum_{\gamma \in \mathcal{F}(\eta)} \prod_{(x, y) \in \gamma} |\omega|(x, y), \quad x \in \eta. \quad (5.11)$$

Here $\mathcal{F}(\eta)$ is the set of trees with the set of vertices η .

4 Construction of limiting Gibbs processes

Theorem 5.4 Let Φ be a \mathcal{P} -stable pair potential in X which is b -regular for a . If $e^{b+a}\rho$ is a Radon measure, then there exists a unique process P_z in X of infinite order having

correlation function r_z , which is the limiting Gibbs process of the sequence $(Q_{z,A_n})_n$ in the weak sense.

The proof of this theorem is based on the following lemma.

Lemma 5.5 ([11]) Let $(P_n)_n$ be a sequence of point processes in X of infinite order satisfying the conditions: for each k the limits $\tilde{v}^k(f) = \lim_{n \rightarrow \infty} \tilde{v}_{P_n}^k(f)$, $f \in \mathcal{H}(X^k)$, exist and $\sum_{\ell=1}^{\infty} v^\ell(A^\ell)^{-\frac{1}{2\ell}} = +\infty$ for each bounded A . Here $v^\ell(A^\ell) = \sum_{\mathcal{J}} \tilde{v}^{|\mathcal{J}|}(A^{|\mathcal{J}|})$, where the summation is over all partitions of $\{1, \dots, \ell\}$ into non-empty subsets. Then there exists one and only one point process P in X of infinite order such that $P_n \Rightarrow P$ and $\tilde{v}_P^k = \tilde{v}^k$ for each k .

Lemma 5.5 combined with the Ruelle bound completes the proof of Theorem 5.4. We consider below the case where $z(x) \equiv z > 0$, $x \in X$.

Proposition 5.6 Under the conditions of Theorem 5.4, $r_z(\xi) = z^{|\xi|} \int_{\mathcal{X}} G(\xi, \eta) \Lambda_{\rho_z}(d\eta)$ satisfies the Kirkwood-Salsburg (K-S) equation:

$$(K\Sigma_{z\rho}) \quad r_z(\xi) = z e^{-W_\Phi(x, \xi)} \cdot \int_{\mathcal{X}} K(x, \eta) r_z(\xi_x + \eta) \Lambda_\rho(d\eta), \quad x \in \xi \neq o,$$

where $K(x, \eta) := \prod_{y \in \eta} \omega(x, y)$.

The proof follows from Theorem 5.2, Minlos' formula and the fact that the Ursell kernel satisfies the equations ([PU09], [Ru69]): $G(o, \eta) = \delta_{o, \eta}$ and

$$G(\xi, \eta) = e^{-W_\Phi(x, \xi_x)} \int_{\mathcal{X}} K(x, \nu) G(\xi_x + \nu, \eta - \nu) \Lambda'_\eta(d\nu), \quad x \in \xi \neq o.$$

Theorem 5.7 Let Φ be a \mathcal{P} -stable b -regular potential for a . If $e^{b+a\rho}$ is a Radon measure and $\sup_x a(x) = C < \infty$ and if the activity satisfies $0 < z < (eC)^{-1}$, the $(K\Sigma_{z\rho})$ equation has a unique solution and the correlation function r_z of the process P_z is this unique solution.

Proof. We follow [10] and [3]. Let \mathcal{E}_δ , $\delta > 0$, be the Banach space of all complex valued measurable functions $\varphi : \mathfrak{X}_+ \rightarrow \mathbb{C}$ such that

$$\|\varphi\|_\delta = \sup_{\xi \in \mathfrak{X}_+} \frac{|\varphi(\xi)|}{\delta^{|\xi|} \prod_{x \in \xi} e^{(a+b)(x)}} < +\infty, \quad (5.12)$$

where $|\xi| = \xi(X)$ denotes the number of particles in ξ . Since r_z satisfies the *Ruelle bound* (5.9), the correlation function r_z belongs to \mathcal{E}_δ with the norm ≤ 1 if $z \leq \delta$.

We define on \mathcal{E}_δ the linear operator K by

$$K\varphi(\varepsilon_x) = z \int_{\mathfrak{X}_+} K(x, \eta) \varphi(\eta) \Lambda_\rho(d\eta), \quad x \in X, \quad (5.13)$$

$$K\varphi(\xi) = z e^{-W_\Phi(x, \xi)} \cdot \int_{\mathfrak{X}} K(x, \eta) \varphi(\eta + \xi_x) \Lambda_\rho(d\eta), \quad x \in \xi \neq o. \quad (5.14)$$

Using the operator K , we can write the K-S equation as an integral equation in the Banach space \mathcal{E}_δ : $r_z = Kr_z + \alpha_z$, where $\alpha_z(\xi) = 0$ if $\xi(X) > 1$ and $\alpha_z(\varepsilon_x) = z$. For sufficiently small $z > 0$ the operator K is bounded. Indeed let $\varphi \in \mathcal{E}_\delta$ with $\|\varphi\| \leq 1$. Then by \mathcal{P} -stability and b -regularity of Φ for every $x \in \xi \in \mathfrak{X}$,

$$\begin{aligned} |(K\varphi)|(\xi) &\leq z e^{b(x)} \int_{\mathfrak{X}} |\omega_x|(\eta) \delta^{|\eta|+|\xi|-1} e^{(\eta+\xi_x)(b+a)} \Lambda_\rho(d\eta) \\ &\leq z \delta^{|\xi|-1} e^{\xi(b+a)} \cdot \exp\left(\delta \rho(|\omega_x| e^{b+a})\right) \leq \frac{z e^{\delta C}}{\delta} \delta^{|\xi|} e^{\xi(b+a)}. \end{aligned}$$

Thus, if the parameters z and δ satisfy the condition $z e^{\delta C} \delta^{-1} < 1$, then $\|K\|_\delta < 1$ and the K-S equation has a unique solution. In particular, if we take $\delta = \frac{1}{C}$, this condition on z becomes $0 < z < (eC)^{-1}$. A more detailed discussion of the choice of δ can be found in [3]. \square

5 Uniqueness of Gibbs processes

In a final step we show that Gibbs processes G for Φ with activity z have correlation functions which solve the K-S equation in the same range of the parameter z . This implies that the Gibbs process G , if it exists, coincides with P_z .

We use the notion of Gibbs process introduced in [6] as a solution of an integration-by-parts formula. A point processes G is called a *Gibbs process for* (Φ, ρ) , if for all $h \in F_+$

$$(\Sigma_\rho) \int_{\mathcal{M}} \int_X h(x, \mu) \mu(dx) G(d\mu) = \int_{\mathcal{M}} \int_X h(x, \mu + \varepsilon_x) \exp(-W_\Phi(x, \mu)) \rho(dx) G(d\mu).$$

We then write $G \in \mathcal{G}(\Phi, \rho)$. This is equivalent to saying that G is a Gibbs process for (Φ, ρ) in the sense of Dobrushin, Lanford and Ruelle, cf. [6].

From now on we assume that Φ is \mathcal{P} -stable, modified b -regular for a and $e^{b+a}\rho$ is a Radon measure. Then

Lemma 5.8 ([6]) Every $G \in \mathcal{G}(\Phi, \rho_z)$ is of infinite order and its correlation function is given by

$$g_z(\xi) = z^{|\xi|} \int_{\mathcal{M}^{\cdot}} \exp(-W_{\Phi}(\xi, \mu)) G(d\mu), \quad \xi \in \mathfrak{X}. \quad (5.15)$$

Furthermore, G is uniquely determined by its correlation functions.

Note that by modified regularity of Φ the conditional energy $W_{\Phi}(x, \mu)$ can be extended to the whole $\mathcal{M}^{\cdot}(X)$ and remains \mathcal{P} -stable with the same function b . Due to the \mathcal{P} -stability, the Ruelle bound takes the form $g_z(\xi) \leq z^{|\xi|} \prod_{x \in \xi} e^{b(x)}$. Using ideas of Sabine Jansen [1], we then obtain

Proposition 5.9 Let G_z be a Gibbs process in X for (Φ, ρ_z) . Then its correlation function g_z solves the $(K\Sigma_{z\rho})$ equation.

In the view of the Ruelle bound we are in the situation as we had been above for the correlation function r_z .

Lemma 5.10 Let Φ be a \mathcal{P} -stable pair potential satisfying above mentioned conditions. Assume also $0 < z < (eC)^{-1}$. Then g_z coincides with r_z which implies that G_z coincides with the limiting Gibbs process P_z .

Thus we arrive at the main result of this paper

Theorem 5.11 For all $0 < z < (eC)^{-1}$ the collection $\mathcal{G}(\Phi, \rho_z)$ of Gibbs processes is either empty or the singleton $\{P_z\}$.

For a large class of hard-core potentials we show in [9] that indeed for all $0 < z < (eC)^{-1}$ the set $\mathcal{G}(\Phi, \rho_z)$ is not empty and therefore reduced to a unique element constructed as the limiting Gibbs process.

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