

On Existence of Solutions for Some Hyperbolic-Parabolic Type Chemotaxis Systems*

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Abstract: In this paper, we discuss the local and global existence of weak solutions for some hyperbolic-parabolic systems modelling chemotaxis.

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1 Introduction

The earliest model for chemosensitive movement has been developed by Keller and Segel [1,2,3], which we call it as KS model. Assume that in absence of any external signal the spread of a population $u(t, x)$ is described by the diffusion equation

$$u_t = d\Delta u, \tag{1}$$

where $d > 0$ is the diffusion constant. We define the net flux as $j = -d\nabla u$. If there is some external signal s , we just assume that it results in a chemotactic velocity β . Then the flux is

$$j = -d\nabla u + \beta u. \tag{2}$$

To be more specific, we assume that the chemotactic velocity β has the direction of the gradient ∇s and that the sensitivity χ to the gradient depends on the signal concentration $s(t, x)$, then $\beta = \chi(s)\nabla s$.

We use this modified flux in (2) to obtain the parabolic chemotaxis equation

$$u_t = \nabla(d\nabla u - \chi(s)\nabla s \cdot u). \tag{3}$$

If $\chi(s)$ is positive, which means that the chemotactic velocity is in direction of s , we call it positive bias, whereas $\chi < 0$ is called negative bias.

To our general knowledge, the external signal is produced by the individuals and decays, which is described by a nonlinear function $g(s, u)$. We assume that the spatial spread of the external signal is driven by diffusion. Then the full system for u and s reads

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$$u_t = \nabla(d\nabla u - \chi(s)\nabla s \cdot u), \quad (4)$$

$$\tau s_t = d\Delta s + g(s, u), \quad (5)$$

the time constant $0 \leq \tau \leq 1$ indicates that the spatial spread of the organisms u and the signal s are on different time scales. The case $\tau = 0$ corresponds to a quasi-steady state assumption for the signal distribution. When we assume that the spatial spread of external signal is driven by wave motion, then the equation (5) would be replaced by

$$s_{tt} = d\Delta s + g(s, u). \quad (6)$$

The full system for u and s becomes

$$u_t = \nabla(d\nabla u - \chi(s)\nabla s \cdot u), \quad (7)$$

$$s_{tt} = d\Delta s + g(s, u), \quad (8)$$

which is called as hyperbolic-parabolic chemotaxis system.

2 Main Results

Let us consider the following problem:

$$\begin{aligned} u_t &= \nabla(\nabla u - \chi u \nabla v) & \text{in } (0, T) \times \Omega, \\ v_{tt} &= \Delta v + g(u, v) & \text{in } (0, T) \times \Omega, \\ \frac{\partial u}{\partial n} &= \frac{\partial v}{\partial n} = 0, & \text{on } (0, T) \times \partial\Omega, \end{aligned} \quad (9)$$

with initial data

$$u(0, \cdot) = u_0, \quad v(0, \cdot) = \varphi, \quad v_t(0, \cdot) = \psi \quad \text{in } \Omega,$$

where $\Omega \subset \mathbf{R}^n$, a bounded open domain with smooth boundary $\partial\Omega$, χ is a nonnegative constant.

Choose a constant σ , which satisfies

$$1 < \sigma < 2 \quad (10)$$

and

$$n < 2\sigma < n + 2 \quad (11)$$

It is easy to check that (10) and (11) can be simultaneously satisfied in the case of $1 \leq n \leq 3$.

Our main results are

Theorem 4.1. *Under the conditions (10) and (11), if $g(u, v) = -\gamma v + f(u)$ and $f \in C^2(\mathbf{R})$, then for each initial data $u_0 \in H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$, $\varphi \in H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}$, $\psi \in H^1(\Omega)$, the problem (9) has a unique local solution $(u, v) \in X_{t_0} \times Y_{t_0}$ for some $t_0 > 0$.*

Theorem 5.1. Let $n = 1$ and $\sigma = \frac{5}{4}$, if $g(u, v) = -\gamma v + f(u)$ and $f \in C_0^2(\mathbf{R})$, then for each initial data $u_0 \in H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$ and $u_0 \geq 0$, $\varphi \in H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}$ and $\psi \in H^1(\Omega)$, the problem (9) has a unique global solution $(u, v) \in X_\infty \times Y_\infty$.

Where we define

$$X_{t_0} = C([0, t_0], H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\})$$

$$Y_{t_0} = C([0, t_0], H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}) \cap C^1([0, t_0], H^1(\Omega))$$

3 Some Basic Lemmas

For $g(u, v) = -\gamma v + f(u)$, and γ is a constant, $f(x) \in C^2(\mathbf{R})$. We divide the system (9) into two parts:

$$\begin{cases} u_t = \nabla(\nabla u - \chi u \nabla v) & \text{in } (0, T) \times \Omega \\ \frac{\partial u}{\partial n} = 0 & \text{on } (0, T) \times \partial\Omega \\ u(0, \cdot) = u_0 & \text{in } \Omega, \end{cases} \quad (12)$$

and

$$\begin{cases} v_{tt} = \Delta v - \gamma v + f(u) & \text{in } (0, T) \times \Omega \\ \frac{\partial v}{\partial n} = 0 & \text{on } (0, T) \times \partial\Omega \\ v(0, \cdot) = \varphi, \quad v_t(0, \cdot) = \psi & \text{in } \Omega. \end{cases} \quad (13)$$

We have

Lemma 3.1. For any $T > 0$, and

$$\varphi \in H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}, \quad \psi \in H^1(\Omega), \quad f(u(t, \cdot)) \in C([0, T]; H^1(\Omega)),$$

then (13) has a unique solution v , satisfying

$$v \in C([0, T]; H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}), \quad v_t \in C([0, T]; H^1(\Omega)), \quad v_{tt} \in C([0, T]; L^2(\Omega)),$$

and

$$\|v(t, \cdot)\|_{H^2(\Omega)} + \|v_t(t, \cdot)\|_{H^1(\Omega)} \leq e^{cT} (\|\varphi\|_{H^2(\Omega)} + \|\psi\|_{H^1(\Omega)} + \int_0^T \|f(u(\tau, \cdot))\|_{H^1(\Omega)} d\tau), \quad \forall t \in [0, T], \quad (14)$$

where $c > 0$ is a constant which is independent of T .

Proof: Set $v_t = w$, we have following system

$$\begin{cases} v_t = w, \\ w_t = \Delta v - \gamma v + f(u). \end{cases} \quad (15)$$

Thus we can write it in a abstract form:

$$\begin{cases} U_t = LU + F(U) & \text{in } X = H^1(\Omega) \times L^2(\Omega), \\ U_0 = U(0, x) = (\varphi, \psi), \end{cases} \quad (16)$$

where $L(v, w) = (w, \Delta v - v)$ for $(v, w) \in D(L)$, $D(L) = H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\} \times H^1(\Omega)$ and $F(v, w) = (0, (1 - \gamma)v + f(u))$.

Define the inner product in X as

$$\langle (v, w), (v', w') \rangle_X = (v, v')_{H^1} + (w, w')_{L^2},$$

where $(\cdot, \cdot)_{H^1}$ and $(\cdot, \cdot)_{L^2}$ represent the inner products in H^1 and L^2 respectively, then X is a Hilbert space.

For $U = (v, w) \in D(L)$, we have

$$\begin{aligned} \langle LU, U \rangle_X &= \langle (w, \Delta v - v), (v, w) \rangle_X \\ &= (w, v)_{H^1} + (\Delta v - v, w)_{L^2} \\ &= (w, v)_{H^1} + (\Delta v, w)_{L^2} - (v, w)_{L^2} \\ &= (w, v)_{H^1} - (\nabla v, \nabla w)_{L^2} - (v, w)_{L^2} \\ &= 0 \end{aligned} \tag{17}$$

Otherwise, for $U = (v, w) \in D(L)$, $U' = (v', w') \in X$,

$$\begin{aligned} \langle L(v, w), (v', w') \rangle_X &= \langle (w, \Delta v - v), (v', w') \rangle_X \\ &= (w, v')_{H^1} + (\Delta v - v, w')_{L^2} \\ &= (w, v')_{H^1} + (\Delta v, w')_{L^2} - (v, w')_{L^2} \end{aligned} \tag{18}$$

If $\langle L(v, w), (v', w') \rangle_X$ is bounded for each $(v, w) \in D(L)$, then $(w, v')_{H^1}$, $(\Delta v, w')_{L^2}$ and $(v, w')_{L^2}$ are bounded for each $(v, w) \in D(L)$, which means that

$$v' \in H^2 \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}, \quad w' \in H^1, \tag{19}$$

that implies $D(L^*) \subset D(L)$. On the other hand, from (17) and the lemma in [6, p9], we know that

$$L^* = -L.$$

Thus we know that L is a generator of a unitary operator group. It is easy to check that for $f(u(t, \cdot)) \in C([0, T], H^1(\Omega))$,

$$F : X \rightarrow X,$$

and

$$\|F(U_1) - F(U_2)\|_X \leq c \|U_1 - U_2\|_X \quad U_i \in X, \quad i = 1, 2,$$

where $\|(v, w)\|_X^2 = \|v\|_{H^1}^2 + \|w\|_{L^2}^2$.

Now we can declare that (16) has a unique solution

$$U \in C^1([0, T], X) \cap C([0, T], D(L)) \quad \text{for each } U_0 \in D(L), \tag{20}$$

which means that for each $(\varphi, \psi) \in D(L)$, (13) has a unique solution

$$v \in C([0, T], H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}), \quad v_t \in C([0, T], H^1(\Omega)) \quad \text{and} \quad v_{tt} \in C([0, T], L^2(\Omega)).$$

Next, we estimate the norm of v . By using the semigroup notation $T(t) = e^{tL}$, we have

$$U = T(t)U_0 + \int_0^t T(t-s)F(U)ds. \tag{21}$$

Since $L = -L^*$, and in terms of (17), we have that

$$\langle LU, U \rangle_X = 0 \quad \text{for each } U \in D(L),$$

and

$$\langle L^*U, U \rangle_X = \langle -LU, U \rangle_X = 0 \quad \text{for each } U \in D(L).$$

Hence L generates a strongly continuous contractive semigroup on Hilbert space X (cf. [4, 5]), in other words, we have

$$\|e^{tL}\| = \|T(t)\| \leq 1. \quad (22)$$

So we know that

$$\begin{aligned} \|U(t)\|_{H^2 \times H^1} &\leq \|T(t)U_0\|_{H^2 \times H^1} + \int_0^t \|T(t-s)F(U(s))\|_{H^2 \times H^1} ds \\ &\leq \|U_0\|_{H^2 \times H^1} + \int_0^t \|F(U)\|_{H^2 \times H^1} ds \\ &= \|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^t \|(1-\gamma)v + f(u)\|_{H^1} ds \\ &\leq \|\varphi\|_{H^2} + \|\psi\|_{H^1} + c \int_0^t \|v\|_{H^1} ds + \int_0^t \|f(u)\|_{H^1} ds \\ &\leq \|\varphi\|_{H^2} + \|\psi\|_{H^1} + c \int_0^t \|U\|_{H^2 \times H^1} ds + \int_0^t \|f(u)\|_{H^1} ds, \quad 0 \leq t \leq T. \end{aligned} \quad (23)$$

From Gronwall's inequality, we know that

$$\begin{aligned} \|U\|_{H^2 \times H^1} &\leq e^{ct} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^t \|f(u)\|_{H^1} ds) \quad 0 \leq t \leq T, \\ &\leq e^{cT} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^T \|f(u)\|_{H^1} ds), \end{aligned} \quad (24)$$

which implies the estimate (14) and the uniqueness follows.

If Ω is a bounded open domain with smooth boundary, in which we can consider the Neumann boundary condition. As we know that the $e^{t\Delta}$ defines a holomorphic semigroup on the Hilbert space $L^2(\Omega)$, so we have that

$$f \in L^2(\Omega) \Rightarrow \|e^{t\Delta}f\|_{H^2(\Omega)} \leq \frac{c}{t} \|f\|_{L^2(\Omega)}, \quad (25)$$

where $D(\Delta) = \{u \in H^2(\Omega), \frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$.

Applying interpolation to (25), it yields

$$\|e^{t\Delta}f\|_{H^\sigma(\Omega)} \leq ct^{-\frac{\sigma}{2}} \|f\|_{L^2(\Omega)} \quad \text{for } 0 \leq \sigma \leq 2, 0 < t \leq 1. \quad (26)$$

Take $Y = H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$ and $Z = L^2(\Omega)$, $\Phi(u) = -\chi \nabla v \nabla u - \chi \Delta v \cdot u$. Then For $v \in Y_{t_0}$, and from the lemma in [4, p273], we can declare that

Lemma 3.2. *For each $u_0 \in Y$ and $v \in Y_{t_0}$, σ and n satisfy the conditions (10) and (11), then the problem (12) has a unique solution*

$$u \in X_{t_0} = C([0, t_0], H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}).$$

Proof: If we can show that $\Phi : Y \rightarrow Z$ is a locally Lipschitz map, then the lemma 3.2 is true. In fact, for arbitrary $u_1, u_2 \in Y$ and $v \in Y_{t_0}$, the difference

$$\Phi(u_1) - \Phi(u_2) = -\chi \nabla v \nabla (u_1 - u_2) - \chi \Delta v \cdot (u_1 - u_2).$$

That is

$$\begin{aligned} \|\Phi(u_1) - \Phi(u_2)\|_Z &= \|\Phi(u_1) - \Phi(u_2)\|_{L^2} \\ &\leq \|\chi \nabla v \nabla(u_1 - u_2)\|_{L^2} + \|\chi \Delta v \cdot (u_1 - u_2)\|_{L^2}. \end{aligned}$$

By Sobolev imbedding theorems, we have

$$H^1(\Omega) \hookrightarrow L^\infty(\Omega), \text{ for } n = 1,$$

$$H^1(\Omega) \hookrightarrow L^q(\Omega), \quad 1 < q < \infty, \text{ for } n = 2,$$

$$H^1(\Omega) \hookrightarrow L^{\frac{2n}{n-2}}(\Omega), \text{ for } n = 3.$$

Thus in terms of (10) and (11), we know that $H^1(\Omega) \hookrightarrow L^{\frac{n}{\sigma-1}}(\Omega)$ and $H^{\sigma-1}(\Omega) \hookrightarrow L^{\frac{2n}{n-2(\sigma-1)}}(\Omega)$ for $n = 2, 3$.

Firstly we estimate $\|\chi \nabla v \nabla(u_1 - u_2)\|_{L^2}$. If $n = 1$, then

$$\begin{aligned} &\|\chi \nabla v \nabla(u_1 - u_2)\|_{L^2} \\ &\leq \chi \|\nabla(u_1 - u_2)\|_{L^2} \|\nabla v\|_{L^\infty} \\ &\leq c \|u_1 - u_2\|_{H^1} \|\nabla v\|_{H^1} \\ &\leq c \|u_1 - u_2\|_{H^\sigma} \|v\|_{H^2}. \end{aligned}$$

If $n = 2, 3$, then

$$\begin{aligned} &\|\chi \nabla v \nabla(u_1 - u_2)\|_{L^2} \\ &\leq \chi \|\nabla(u_1 - u_2)\|_{L^{\frac{2n}{n-2(\sigma-1)}}} \|\nabla v\|_{L^{\frac{n}{\sigma-1}}} \\ &\leq c \|u_1 - u_2\|_{H^\sigma} \|v\|_{H^2}. \end{aligned}$$

Hence for $n = 1, 2, 3$, we have that

$$\|\chi \nabla v \nabla(u_1 - u_2)\|_{L^2} \leq c \|u_1 - u_2\|_{H^\sigma} \|v\|_{H^2}.$$

Similarly, we have

$$\begin{aligned} &\|\chi \Delta v \cdot (u_1 - u_2)\|_{L^2} \\ &\leq c \|v\|_{H^2} \|u_1 - u_2\|_{L^\infty} \\ &\leq c \|u_1 - u_2\|_{H^\sigma} \|v\|_{H^2}. \end{aligned}$$

Thus we have proved that

$$\|\Phi(u_1) - \Phi(u_2)\|_Z \leq c \|u_1 - u_2\|_Y \|v\|_{H^2},$$

as required.

Lemma 3.3. *Under the conditions (10) and (11), if $u \in X_{t_0}$ is a solution of (12), there exists a constant c which is independent of t_0 , such that*

$$\|u\|_{X_{t_0}} \leq c \|u_0\|_{\sigma,2} + ct_0^{1-\frac{\sigma}{2}} \|v\|_{Y_{t_0}} \cdot \|u\|_{X_{t_0}}, \quad (27)$$

where $\|\cdot\|_{k,p}$ is the norm of Sobolev space $W^{k,p}$.

Proof: Let $T(t) = e^{t\Delta}$, then

$$u(t) = T(t)u_0 - \chi \int_0^t T(t-s) \nabla v \nabla u ds - \chi \int_0^t T(t-s) \Delta v \cdot u ds.$$

By (26), we have $T(t) : L^2(\Omega) \rightarrow H^\sigma(\Omega)$ with norm $c_\sigma t^{-\frac{\sigma}{2}}$. Thus

$$\left\| \int_0^t T(t-s) \nabla v \nabla u ds \right\|_{\sigma,2} \leq c_\sigma t^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t} \|\nabla v(s, \cdot) \nabla u(s, \cdot)\|_2$$

where we use $\|\cdot\|_p$ as the norm of L^p .

By Sobolev imbedding theorem, $H^1(\Omega) \hookrightarrow L^\infty(\Omega)$ for $n = 1$, we have

$$\begin{aligned} \|\nabla v \nabla u\|_2 &\leq \|\nabla v\|_\infty \cdot \|\nabla u\|_2 \\ &\leq c \|v\|_{2,2} \cdot \|u\|_{1,2} \\ &\leq c \|v\|_{2,2} \cdot \|u\|_{\sigma,2}. \end{aligned}$$

For $n = 2, 3$, we have $H^1(\Omega) \hookrightarrow L^{\frac{n}{\sigma-1}}(\Omega)$, $H^{\sigma-1}(\Omega) \hookrightarrow L^{\frac{2n}{n-2(\sigma-1)}}(\Omega)$, thus $f^2 \in L^{\frac{n}{2(\sigma-1)}}$, $g^2 \in L^{\frac{n}{n-2(\sigma-1)}}$ if $f \in H^1$ and $g \in H^{\sigma-1}$. By using Cauchy inequality, we get

$$\|f^2 g^2\|_1 \leq \|f^2\|_{\frac{n}{2(\sigma-1)}} \cdot \|g^2\|_{\frac{n}{n-2(\sigma-1)}}$$

which implies $\|fg\|_2 \leq \|f\|_{\frac{n}{\sigma-1}} \cdot \|g\|_{\frac{2n}{n-2(\sigma-1)}}$. Thus

$$\begin{aligned} \|\nabla v \nabla u\|_2 &\leq \|\nabla v\|_{\frac{n}{\sigma-1}} \cdot \|\nabla u\|_{\frac{2n}{n-2(\sigma-1)}} \\ &\leq c \|\nabla v\|_{1,2} \cdot \|\nabla u\|_{\frac{2n}{n-2(\sigma-1)}} \\ &\leq c \|v\|_{2,2} \cdot \|\nabla u\|_{\sigma-1,2} \leq c \|v\|_{2,2} \cdot \|u\|_{\sigma,2}. \end{aligned}$$

Now we obtain that, for $0 \leq t \leq t_0$,

$$\begin{aligned} \left\| \int_0^t \tau(t-s) \nabla v \nabla u ds \right\|_{\sigma,2} &\leq c_\sigma t^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t} \|\nabla v \nabla u\|_2 \\ &\leq C t^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t} \|v\|_{2,2} \cdot \|u\|_{\sigma,2} \leq C t_0^{1-\frac{\sigma}{2}} \|u\|_{X_{t_0}} \cdot \|v\|_{Y_{t_0}}. \end{aligned}$$

Meanwhile

$$\begin{aligned} &\left\| \int_0^t T(t-s) \Delta v \cdot u \right\|_{\sigma,2} \\ &\leq c_\sigma t^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t} \|\Delta v \cdot u\|_2 \\ &\leq c_\sigma t_0^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t_0} \|u\|_{L^\infty} \cdot \|\Delta v\|_{L^2} \\ &\leq C t_0^{1-\frac{\sigma}{2}} \sup_{0 \leq s \leq t_0} \|u\|_{\sigma,2} \cdot \sup_{0 \leq s \leq t_0} \|v\|_{2,2} \\ &\leq C t_0^{1-\frac{\sigma}{2}} \|u\|_{X_{t_0}} \cdot \|v\|_{Y_{t_0}}. \end{aligned}$$

Finally we can deduce that

$$\begin{aligned} \|u(t)\|_{\sigma,2} &\leq \|T(t)u_0\|_{\sigma,2} + \chi \left\| \int_0^t T(t-s) \nabla v \nabla u ds \right\|_{\sigma,2} \\ &\quad + \chi \left\| \int_0^t T(t-s) \Delta v \cdot u ds \right\|_{\sigma,2} \\ &\leq C \|u_0\|_{\sigma,2} + \chi c c_\sigma t_0^{1-\frac{\sigma}{2}} \|u\|_{X_{t_0}} \cdot \|v\|_{Y_{t_0}}, \quad 0 \leq t \leq t_0, \end{aligned}$$

which implies

$$\|u\|_{X_{t_0}} \leq C \|u_0\|_{\sigma,2} + C t_0^{1-\frac{\sigma}{2}} \|u\|_{X_{t_0}} \|v\|_{Y_{t_0}}.$$

Lemma 3.3 is proved.

4 Local Existence of Solutions

In this section, we establish the local solution of the system (9). Our main result is as follows:

Theorem 4.1. *If σ and n satisfy the conditions (10) and (11), $g(u, v) = -\gamma v + f(u)$ and $f \in C^2(\mathbf{R})$, then for each initial data $u_0 \in H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$, $\varphi \in H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}$, $\psi \in H^1(\Omega)$, the problem (9) has a unique local solution $(u, v) \in X_{t_0} \times Y_{t_0}$ for some $t_0 > 0$.*

Proof: Consider $w \in X_{t_0}$, $w(0, x) = u_0(x)$ and let $v = v(w)$ denote the corresponding solution of the equation:

$$\begin{aligned} v_{tt} &= \Delta v - \gamma v + f(w) \quad \text{in } (0, t_0) \times \Omega, \\ \frac{\partial v}{\partial n} &= 0 \quad \text{on } (0, t_0) \times \partial\Omega, \\ v(0) &= \varphi \quad \text{in } \Omega, \\ v_t(0) &= \psi \quad \text{in } \Omega. \end{aligned} \tag{28}$$

By Lemma 3.1, we have $v \in Y_{t_0}$, and

$$\begin{aligned} \|v(t)\|_{H^2(\Omega)} &\leq e^{c_1 t_0} (\|\varphi\|_{H^2(\Omega)} + \|\psi\|_{H^1(\Omega)} \\ &\quad + \int_0^{t_0} \|f(w(\tau, \cdot))\|_{H^1(\Omega)} d\tau), \quad \forall t \in [0, t_0]. \end{aligned} \tag{29}$$

Secondly, for the solution v of (28), we define $u = u(v(w))$ to be the corresponding solution of

$$\begin{aligned} u_t &= \nabla(\nabla u - \chi u \nabla v) \quad \text{in } (0, t_0) \times \Omega, \\ \frac{\partial u}{\partial n} &= 0 \quad \text{on } (0, t_0) \times \partial\Omega, \\ u(0, x) &= u_0(x) = w(0, x) \quad \text{in } \Omega. \end{aligned} \tag{30}$$

If we define $Gw = u(v(w))$, then Lemma 3.2 shows that

$$G: X_{t_0} \rightarrow X_{t_0}.$$

Take $M = 2c \|u_0\|_{\sigma, 2}$ and a ball

$$B_M = \left\{ w \in X_{t_0} \mid w(0, x) = u_0(x), \|w(t, \cdot)\|_{\sigma, 2} \leq M, 0 \leq t \leq t_0 \right\},$$

where the constant $c \geq 1$ is given by (27). Then we combine the estimates (27) and (29) to obtain

$$\begin{aligned} \|Gw\|_{X_{t_0}} &\leq c \|u_0\|_{\sigma, 2} + ct_0^{1-\frac{\sigma}{2}} \|v\|_{Y_{t_0}} \cdot \|Gw\|_{X_{t_0}} \\ &\leq c \|u_0\|_{\sigma, 2} + ct_0^{1-\frac{\sigma}{2}} e^{c_1 t_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} \\ &\quad + \int_0^{t_0} \|f(w(\tau, \cdot))\|_{H^1} d\tau) \cdot \|Gw\|_{X_{t_0}}. \end{aligned}$$

Since $\|w\|_{1, 2} \leq \|w\|_{\sigma, 2} \leq M$, and $f \in C^2(\mathbf{R})$, we can deduce that

$$\|f(w(\tau, \cdot))\|_{1, 2} \leq \|f\|_{C^2[-M, M]} \cdot M + \|f(0)\|_{L^2},$$

which shows that $\|Gw\|_{X_{t_0}} \leq 2c \|u_0\|_{\sigma, 2}$ for $t_0 > 0$ small enough.

Thus we have proved that, for $t_0 > 0$ small enough, G maps B_M into B_M . Next, we can prove that, for t_0 small enough, G is a contract mapping. In fact, let $w_1, w_2 \in X_u$, and v_1, v_2 denote the corresponding solutions of (28). Then the difference $Gw_1 - Gw_2$ satisfies:

$$\begin{aligned} Gg_1 - Gg_2 &= u_1 - u_2 \\ &= -\chi \int_0^t T(t-s)u_1 \Delta v_1 ds - \chi \int_0^t T(t-s) \nabla u_1 \nabla v_1 ds \\ &\quad + \chi \int_0^t T(t-s)u_2 \nabla v_2 ds + \chi \int_0^t T(t-s) \nabla u_2 \nabla v_2 ds \\ &= -\chi \int_0^t T(t-s)(u_1 \Delta v_1 - u_2 \Delta v_2) ds - \chi \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_2) ds. \end{aligned}$$

Next, we have

$$\begin{aligned} &\left\| \int_0^t T(t-s)(u_1 \Delta v_1 - u_2 \Delta v_2) ds \right\|_{\sigma,2} \\ &\leq \left\| \int_0^t T(t-s)u_1(\Delta v_1 - \Delta v_2) ds \right\|_{\sigma,2} + \left\| \int_0^t T(t-s)(u_1 - u_2) \Delta v_2 ds \right\|_{\sigma,2}. \end{aligned}$$

Since

$$\begin{aligned} &\left\| \int_0^t T(t-s)u_1(\Delta v_1 - \Delta v_2) ds \right\|_{\sigma,2} \\ &\leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|u_1(\Delta v_1 - \Delta v_2)\|_2 \\ &\leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|u_1\|_{L^\infty} \cdot \|\Delta(v_1 - v_2)\|_2 \\ &\leq CMt_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|v_1 - v_2\|_{2,2}, \end{aligned} \tag{31}$$

and

$$\begin{aligned} &\left\| \int_0^t T(t-s)(u_1 - u_2) \Delta v_2 ds \right\|_{\sigma,2} \\ &\leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|(u_1 - u_2) \Delta v_2\|_2 \\ &\leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|v_2\|_{2,2} \cdot \|u_1 - u_2\|_{L^\infty} \\ &\leq ct_0^{1-\frac{\sigma}{2}} \|v_2\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}}. \end{aligned} \tag{32}$$

Thus we have that

$$\begin{aligned} &\left\| \int_0^t T(t-s)(u_1 \Delta v_1 - u_2 \Delta v_2) ds \right\|_{\sigma,2} \\ &\leq Ct_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}} \\ &\quad + Ct_0^{1-\frac{\sigma}{2}} \|v_2\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}}, \quad 0 \leq t \leq t_0. \end{aligned} \tag{33}$$

Similarly, we have

$$\begin{aligned} &\left\| \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_2) ds \right\|_{\sigma,2} \\ &\leq \left\| \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_1) ds \right\|_{\sigma,2} \\ &\quad + \left\| \int_0^t T(t-s)(\nabla u_2 \nabla v_1 - \nabla u_2 \nabla v_2) ds \right\|_{\sigma,2}. \end{aligned}$$

Here

$$\begin{aligned} &\left\| \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_1) ds \right\|_{\sigma,2} \\ &\leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|\nabla v_1 \cdot \nabla(u_1 - u_2)\|_2, \quad 0 \leq t \leq t_0. \end{aligned}$$

As we have done in Lemma 3.3, we can deduce that

$$\begin{aligned} & \left\| \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_1) ds \right\|_{\sigma,2} \\ & \leq Ct_0^{1-\frac{\sigma}{2}} \|v_1\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}}, \quad 0 \leq t \leq t_0. \end{aligned} \quad (34)$$

And we have similarly that

$$\begin{aligned} & \left\| \int_0^t T(t-s)(\nabla u_2 \nabla v_1 - \nabla u_2 \nabla v_2) ds \right\|_{\sigma,2} \\ & \leq ct_0^{1-\frac{\sigma}{2}} \sup_{0 \leq t \leq t_0} \|\nabla u_2 \cdot \nabla(v_1 - v_2)\|_2 \\ & \leq ct_0^{1-\frac{\sigma}{2}} \|u_2\|_{X_{t_0}} \cdot \|v_1 - v_2\|_{Y_{t_0}} \\ & \leq cMt_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}}, \quad 0 \leq t \leq t_0. \end{aligned} \quad (35)$$

Then

$$\begin{aligned} & \left\| \int_0^t T(t-s)(\nabla u_1 \nabla v_1 - \nabla u_2 \nabla v_2) ds \right\|_{\sigma,2} \\ & \leq Ct_0^{1-\frac{\sigma}{2}} \|v_1\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}} + Ct_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}}, \quad 0 \leq t \leq t_0. \end{aligned} \quad (36)$$

Combining the estimates (33) and (36), we have

$$\begin{aligned} & \|Gw_1 - Gw_2\|_{\sigma,2} = \|u_1 - u_2\|_{\sigma,2} \\ & \leq Ct_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}} + Ct_0^{1-\frac{\sigma}{2}} \|v_2\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}} \\ & \quad + Ct_0^{1-\frac{\sigma}{2}} \|v_1\|_{Y_{t_0}} \cdot \|u_1 - u_2\|_{X_{t_0}} + Ct_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}}, \end{aligned}$$

which implies

$$\begin{aligned} & \|Gw_1 - Gw_2\|_{X_{t_0}} \\ & \leq 2Ct_0^{1-\frac{\sigma}{2}} \|v_1 - v_2\|_{Y_{t_0}} + Ct_0^{1-\frac{\sigma}{2}} (\|v_2\|_{Y_{t_0}} + \|v_1\|_{Y_{t_0}}) \cdot \|Gw_1 - Gw_2\|_{X_{t_0}}. \end{aligned}$$

Also, we have

$$\begin{aligned} \|v_1 - v_2\|_{2,2} & \leq e^{c_1 t_0} \int_0^{t_0} \|f(w_1) - f(w_2)\|_{H^1} d\tau \\ & \leq e^{c_1 t_0} \|f\|_{C^2[-M,M]} \int_0^{t_0} \|w_1 - w_2\|_{H^\sigma} d\tau, \end{aligned}$$

and

$$\begin{aligned} \|v_1\|_{2,2} & \leq e^{c_1 t_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^{t_0} \|f(w_1(\tau))\|_{H^1} d\tau) \\ & \leq e^{c_1 t_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + c \int_0^{t_0} (\|w_1(\tau)\|_{H^\sigma} + \|f(0)\|_{H^1}) d\tau) \\ & \leq e^{c_1 t_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + ct_0(M + \|f(0)\|_{L^2})) \\ \|v_2\|_{2,2} & \leq e^{c_1 t_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + ct_0(M + \|f(0)\|_{L^2})). \end{aligned}$$

Thus for $t_0 > 0$ small enough, G is contract.

From process above, we have proved the existence of solution for the problem (9). Since G is contract, then the solution is unique.

5 Global existence of Solutions for $n = 1$

In this section, we establish the global existence and uniqueness of the solution $(u, v) \in X_\infty \times Y_\infty$ of (9) in the case of $n = 1$ and $g(u, v) = -\gamma v + f(u)$. Here we suppose that

$$f(x) \in C_0^2(\mathbf{R}), \quad \sigma = \frac{5}{4}. \quad (37)$$

Observe that, for $n = 1$, $\sigma = \frac{5}{4}$ can simultaneously satisfy the condition (10) and (11). So from the result of Theorem 4.1, the problem (9) has a unique local solution $(u, v) \in X_{t_0} \times Y_{t_0}$ for some $t_0 > 0$ small enough.

Actually we can obtain following more strong result:

Theorem 5.1. *If $n = 1$, $g(u, v) = -\gamma v + f(u)$ and σ and f satisfy the condition (37), then for each initial data $u_0 \in H^\sigma(\Omega) \cap \{\frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$ and $u_0 \geq 0$, $\varphi \in H^2(\Omega) \cap \{\frac{\partial v}{\partial n} = 0 \text{ on } \partial\Omega\}$ and $\psi \in H^1(\Omega)$, the problem (9) has a unique global solution $(u, v) \in X_\infty \times Y_\infty$.*

If $u_0 \geq 0$, then from the first equation of (9), we can deduce that the local solution (u, v) satisfies

$$\|u(t, \cdot)\|_{L^1} = \|u_0\|_{L^1} \quad (38)$$

Next, we have

Lemma 5.2. *Let $s \leq 2$, the local solution $(u, v) \in X_{t_0} \times Y_{t_0}$ of (9), for $g(u, v) = -\gamma v + f(u)$, satisfies*

$$\|v(t, \cdot)\|_{H^s} \leq e^{ct_0}(c_0 + \int_0^{t_0} \|f(u(\tau, \cdot))\|_{H^{s-1}} d\tau), \quad 0 \leq t \leq t_0, \quad (39)$$

where $c_0 = \|\varphi\|_{H^2} + \|\psi\|_{H^1}$ and c is independent of t_0 .

Proof: For $U = (v, w)$ and $F(U) = (0, (1 - \gamma)v + f(u))$, in terms of (21), we know that

$$U = T(t)U_0 + \int_0^t T(t - \tau)F(U(\tau))d\tau$$

where $w = v_t$ and (u, v) is the solution of (9).

By using (22), we know that

$$\begin{aligned} \|U(t)\|_{H^1 \times L^2} &\leq \|T(t)U_0\|_{H^1 \times L^2} + \int_0^t \|T(t - \tau)F(U(\tau))\|_{H^1 \times L^2} d\tau \\ &\leq \|U_0\|_{H^1 \times L^2} + \int_0^t \|F(U(\tau))\|_{H^1 \times L^2} d\tau \\ &= \|\varphi\|_{H^1} + \|\psi\|_{L^2} + \int_0^t \|(1 - \gamma)v + f(u)\|_{L^2} d\tau \\ &\leq \|\varphi\|_{H^1} + \|\psi\|_{L^2} + c \int_0^t \|v\|_{L^2} d\tau + \int_0^t \|f(u)\|_{L^2} d\tau \\ &\leq \|\varphi\|_{H^1} + \|\psi\|_{L^2} + c \int_0^t \|U(\tau)\|_{H^1 \times L^2} d\tau + \int_0^{t_0} \|f(u)\|_{L^2} d\tau, \quad 0 \leq t \leq t_0. \end{aligned} \quad (40)$$

So the Gronwall's inequality indicates

$$\begin{aligned} \|U(t)\|_{H^1 \times L^2} &\leq e^{ct}(\|\varphi\|_{H^1} + \|\psi\|_{L^2} + \int_0^{t_0} \|f(u)\|_{L^2} d\tau) \\ &\leq e^{ct_0}(\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^{t_0} \|f(u)\|_{L^2} d\tau), \quad 0 \leq t \leq t_0. \end{aligned} \quad (41)$$

Since $H^s \times H^{s-1} \subset H^1 \times L^2$ for $s > 1$, we denote $T(t)|_{H^s \times H^{s-1}}$ as the restriction of $T(t)$ on $H^s \times H^{s-1}$, the norm of $T(t)|_{H^s \times H^{s-1}}$ satisfies also the estimate (22). Thus, by similar process of (40) and (41), we can deduce that

$$\|U(t)\|_{H^s \times H^{s-1}} \leq e^{ct_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^{t_0} \|f(u)\|_{H^{s-1}} d\tau), \quad 0 \leq t \leq t_0. \quad (42)$$

If $s < 1$, then $H^1 \times L^2 \subset H^s \times H^{s-1}$, we use Hahn-Banach theorem to get that the operator $T(t)$ can be continuously extended on $H^s \times H^{s-1}$ and the norm of $T(t)$ is invariable. Thus for $s < 1$, we have also that

$$\|U(t)\|_{H^s \times H^{s-1}} \leq e^{ct_0} (\|\varphi\|_{H^2} + \|\psi\|_{H^1} + \int_0^{t_0} \|f(u)\|_{H^{s-1}} d\tau), \quad 0 \leq t \leq t_0. \quad (43)$$

Lemma 5.2 can be deduced directly by (41), (42) and (43).

Proof of theorem 5.1:

For the unique local solution $(u, v) \in X_{t_0} \times Y_{t_0}$ of (9), if we take $s=1/2$ in (39), then

$$\|v(t, \cdot)\|_{H^{\frac{1}{2}}}^2 \leq ce^{t_0} (c_0 + \int_0^{t_0} \|f(u(\tau, \cdot))\|_{H^{-\frac{1}{2}}}^2 d\tau), \quad 0 \leq t \leq t_0. \quad (44)$$

Since $n = 1$, then from Sobolev imbedding theorems, we can deduce that $W^{0,1}(\Omega) \hookrightarrow H^{-\frac{1}{2}}(\Omega)$. Hence we have

$$\begin{aligned} \|v(t, \cdot)\|_{H^{\frac{1}{2}}}^2 &\leq ce^{t_0} (c_0 + \int_0^{t_0} \|f(u(\tau, \cdot))\|_{H^{-\frac{1}{2}}}^2 d\tau) \\ &\leq ce^{t_0} (c_0 + \int_0^{t_0} \|f(u(\tau, \cdot))\|_{L^1}^2 d\tau) \\ &\leq ce^{t_0} (c_0 + \int_0^{t_0} (M_1 \|u\|_{L^1} + \|f(0)\|_{L^1})^2 d\tau) \\ &= ce^{t_0} (c_0 + t_0 (M_1 \|u_0\|_{L^1} + \|f(0)\|_{L^1})^2), \quad 0 \leq t \leq t_0, \end{aligned} \quad (45)$$

where $M_1 = \|f\|_{C^2}$.

On the other hand, for each $s \leq \sigma$ and $0 \leq \sigma_0 < 2$, we have that

$$\begin{aligned} \|u(t, \cdot)\|_{H^s} &\leq c \|u_0\|_{H^\sigma} + ct_0^{1-\frac{\sigma_0}{2}} \|\nabla(u\nabla v)\|_{H^{s-\sigma_0}} \\ &\leq c \|u_0\|_{H^\sigma} + ct_0^{1-\frac{\sigma_0}{2}} \|u\nabla v\|_{H^{s-\sigma_0+1}}, \quad 0 \leq t \leq t_0, \end{aligned} \quad (46)$$

Especially for $s = -\frac{1}{2} + \frac{1}{4}$ and $\sigma_0 = 2 - \frac{1}{8}$, we have

$$\|u(t, \cdot)\|_{H^{-\frac{1}{2} + \frac{1}{4}}} \leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u\nabla v\|_{H^{-1-\frac{1}{8}}}, \quad 0 \leq t \leq t_0. \quad (47)$$

By Sobolev imbedding theorems and (45),

$$\begin{aligned} \|u\nabla v\|_{H^{-1-\frac{1}{8}}} &\leq c \|u\|_{H^{-1-\frac{1}{8}}} \cdot \|\nabla v\|_{W^{-1-\frac{1}{8}, \infty}} \\ &\leq c \|u\|_{H^{-1}} \cdot \|\nabla v\|_{H^{-\frac{1}{2}}} \\ &\leq c \|u\|_{L^1} \cdot \|v\|_{H^{\frac{1}{2}}} \\ &\leq c \|u_0\|_{L^1} \cdot e^{\frac{1}{2}t_0} (c_0^{\frac{1}{2}} + t_0^{\frac{1}{2}} (M_1 \|u_0\|_{L^1} + \|f(0)\|_{L^1})), \quad 0 \leq t \leq t_0. \end{aligned} \quad (48)$$

Thus

$$\begin{aligned} \|u(t, \cdot)\|_{H^{-\frac{1}{4}}} &\leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u\nabla v\|_{H^{-1-\frac{1}{8}}} \\ &\leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u_0\|_{L^1} \cdot e^{\frac{1}{2}t_0} (c_0^{\frac{1}{2}} + t_0^{\frac{1}{2}} (M_1 \|u_0\|_{L^1} + \|f(0)\|_{L^1})), \quad 0 \leq t \leq t_0. \end{aligned} \quad (49)$$

Take $s = \frac{1}{2} + \frac{1}{4} = \frac{3}{4}$ in (39), then (39) and (49) give

$$\begin{aligned}
& \|v(t, \cdot)\|_{H^{\frac{3}{4}}}^2 \leq ce^{t_0} (c_0 + \int_0^{t_0} \|f(u(\tau, \cdot))\|_{H^{\frac{3}{4}-1}}^2 d\tau) \\
& \leq ce^{t_0} (c_0 + t_0 (M_1 \sup_{0 \leq \tau \leq t_0} \|u(\tau, \cdot)\|_{H^{-\frac{1}{4}}} + \|f(0)\|_{H^{-\frac{1}{4}}})^2) \\
& \leq ce^{t_0} (c_0 + t_0 (M_1 (c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u_0\|_{L^1} \cdot e^{\frac{1}{2}t_0} (c_0^{\frac{1}{2}} \\
& + t_0^{\frac{1}{2}} (M_1 \|u_0\|_{L^1} + \|f(0)\|_{L^1})) + \|f(0)\|_{H^{-\frac{1}{4}}}))^2), \quad 0 \leq t \leq t_0.
\end{aligned} \tag{50}$$

Take $s = -\frac{1}{2} + \frac{1}{4} + \frac{1}{4} = 0$ and $\sigma_0 = 2 - \frac{1}{8}$ in (46) again, we obtain that

$$\begin{aligned}
& \|u(t, \cdot)\|_{L^2} \leq c \|u_0\|_{H^\sigma} + ct_0^{1-\frac{\sigma_0}{2}} \|\nabla(u\nabla v)\|_{H^{-\sigma_0}} \\
& \leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u\nabla v\|_{H^{-\sigma_0+1}} \\
& \leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u\nabla v\|_{H^{-1+\frac{1}{8}}}, \quad 0 \leq t \leq t_0.
\end{aligned} \tag{51}$$

Since we know that

$$\begin{aligned}
& \|u\nabla v\|_{H^{-1+\frac{1}{8}}} \leq c \|u\|_{H^{-1+\frac{1}{8}}} \cdot \|\nabla v\|_{W^{-1+\frac{1}{8}, \infty}} \\
& \leq c \|u\|_{H^{-\frac{1}{4}}} \cdot \|\nabla v\|_{H^{-\frac{1}{2}+\frac{1}{4}}} \\
& \leq c \|u\|_{H^{-\frac{1}{4}}} \cdot \|v\|_{H^{\frac{3}{4}}}, \quad 0 \leq t \leq t_0.
\end{aligned} \tag{52}$$

We can get that

$$\begin{aligned}
& \|u(t, \cdot)\|_{L^2} \leq c \|u_0\|_{H^\sigma} + ct_0^{1-\frac{\sigma_0}{2}} \|\nabla(u\nabla v)\|_{H^{-\sigma_0}} \\
& \leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \|u\nabla v\|_{H^{-1+\frac{1}{8}}} \\
& \leq c \|u_0\|_{H^\sigma} + ct_0^{\frac{1}{16}} \cdot \|u\|_{H^{-\frac{1}{4}}} \cdot \|v\|_{H^{\frac{3}{4}}}, \quad 0 \leq t \leq t_0.
\end{aligned} \tag{53}$$

From (49) and (50), we have obtained that $\|u(t, \cdot)\|_{L^2}$ grows by a bounded manner in time.

Again we take $s = \frac{1}{2} + \frac{1}{4} + \frac{1}{4} = 1$ in (39), then (39) and (53) imply that $\|v(t, \cdot)\|_{H^1}$ grows also by a bounded manner in time.

Taking $s = -\frac{1}{2} + \frac{1}{4} + \frac{1}{4} + \frac{1}{4} = \frac{1}{4}$ and $\sigma_0 = 2 - \frac{1}{8}$ in (46) once more, since $\|v(t, \cdot)\|_{H^1}$ grows by a bounded manner in time, similar to which we have done in (51), (52) and (53), we can deduce that $\|u(t, \cdot)\|_{H^{\frac{1}{4}}}$ grows by a bounded manner in time.

Let us repeat processes above four times, we can prove that $\|u(t, \cdot)\|_{H^{\frac{5}{4}}}$ and $\|v(t, \cdot)\|_{H^2}$ grow by a bounded manner in time, as required.

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