Universität Potsdam


Giovanni Conforti | Sylvie Rœlly
Reciprocal class of random walks on an Abelian group

Preprints des Instituts für Mathematik der Universität Potsdam 4 (2015) I

Preprints des Instituts für Mathematik der Universität Potsdam 4 (2015) 1

Giovanni Conforti | Sylvie Roelly

## Reciprocal class of random walks on an Abelian group

## Bibliografische Information der Deutschen Nationalbibliothek

Die Deutsche Nationalbibliothek verzeichnet diese Publikation in der Deutschen Nationalbibliografie; detaillierte bibliografische Daten sind im Internet über http://dnb.dnb.de abrufbar.

## Universitätsverlag Potsdam 2015

http://verlag.ub.uni-potsdam.de/
Am Neuen Palais 10, 14469 Potsdam
Tel.: +49 (0)331 9772533 / Fax: 2292
E-Mail: verlag@uni-potsdam.de
Die Schriftenreihe Preprints des Instituts für Mathematik der Universität Potsdam wird herausgegeben vom Institut für Mathematik der Universität Potsdam.

ISSN (online) 2193-6943
Kontakt:
Institut für Mathematik
Am Neuen Palais 10
14469 Potsdam
Tel.: +49 (0)331 9771028
WWW: http://www.math.uni-potsdam.de
Titelabbildungen:

1. Karla Fritze | Institutsgebäude auf dem Campus Neues Palais
2. Nicolas Curien, Wendelin Werner | Random hyperbolic triangulation

Published at: http://arxiv.org/abs/1105.5089
Das Manuskript ist urheberrechtlich geschützt.
Online veröffentlicht auf dem Publikationsserver der Universität Potsdam URL http://pub.ub.uni-potsdam.de/volltexte/2015/7260/
URN urn:nbn:de:kobv:517-opus-72604
http://nbn-resolving.de/urn:nbn:de:kobv:517-opus-72604

# Reciprocal class of random walks on an Abelian group 

GIOVANNI CONFORTI ${ }^{1,2, *}$ and SYLVIE ROELLY ${ }^{2, * *}$<br>${ }^{1}$ Università degli Studi di Padova, Dipartimento di Matematica Pura e Applicata. Via Trieste 63, 35121 Padova, Italy. E-mail: * giovanniconfort@gmail.com<br>${ }^{2}$ Institut für Mathematik der Universität Potsdam, Am Neuen Palais 10, 14469 Potsdam, Germany. E-mail: ${ }^{* *}$ roelly@math.uni-potsdam.de (Corresponding author)


#### Abstract

Processes having the same bridges as a given reference Markov process constitute its reciprocal class. In this paper we study the reciprocal class of a continuous time random walk with values in a countable Abelian group, we compute explicitly its reciprocal characteristics and we present an integral characterization of it. Our main tool is a new iterated version of the celebrated Mecke's formula from the point process theory, which allows us to study, as transformation on the path space, the addition of random loops. Thanks to the lattice structure of the set of loops, we even obtain a sharp characterization. At the end, we discuss several examples to illustrate the richness of reciprocal classes. We observe how their structure depends on the algebraic properties of the underlying group.


Keywords: Reciprocal class, stochastic bridge, random walk on Abelian group.

## Contents

Introduction ..... 2
1 Random walk on Abelian groups ..... 3
1.1 The random walk as Poisson random measure ..... 3
1.2 An integral characterization ..... 4
1.3 An iterated formula ..... 5
2 The reciprocal class and its characterization ..... 7
2.1 Bridges of the random walk and their mixtures ..... 7
2.2 Loops and their skeletons ..... 7
2.3 Main result: an integral characterization of the reciprocal class ..... 8
2.4 Proof of the main theorem ..... 9
3 Examples ..... 16
3.1 The group $G=\mathbb{Z}$ is infinite ..... 16
3.1.1 The finite support $G_{\nu}$ of the jump measure $\nu$ contains $\{-1,1\}$. ..... 16
3.1.2 $\quad G_{\nu}=\{1,2\}$. ..... 18
$3.2 G$ is the cyclic group $\mathbb{Z} / N \mathbb{Z}$ ..... 18
3.2.1 The support $G_{\nu}$ of the jump measure reduces to $\{\mathbf{- 1}, \mathbf{1}\}$. ..... 18
3.2.2 The support $G_{\nu}$ of the jump measure covers $\mathbb{Z} / N \mathbb{Z} \backslash\{\mathbf{0}\}$. ..... 19
3.3 The state space is a product group ..... 19
Acknowledgments ..... 20
References . ..... 21

## Introduction

Given a reference Markov probability on the time interval [ 0,1 ], its reciprocal class is the set of all probability measures which can be written as a mixture of its bridges. All the elements of the class enjoy the reciprocal property, which is a weaker version of the Markov property since it is a (time) Markov field property. For a detailed comparison between the Markov property and the reciprocal one, we refer to the recent survey [10]. Reciprocal probabilities were first introduced by Schrödinger in [18] to study the dynamics of a Brownian particle with prescribed laws at the initial and the final times, see e.g. [9]. Bernstein then highlighted in [1] the importance of non Markovian reciprocal processes. Jamison initiated later in a series of papers $[6],[7],[8]$ a rigorous mathematical study of reciprocal processes, by partitioning them into classes and underlining the importance of invariants called reciprocal characteristics. The problem of computing, interpreting reciprocal characteristics and then using them to characterize reciprocal classes has attracted the attention of many authors in the context of diffusions (see e.g. [2],[19],[20], [16]).
The study of reciprocal classes of particular jump processes has been started recently by Murr with the case of counting processes, see [12] and [5]. Then results concerning the characterization of the reciprocal classes of a compound Poisson process have been obtained in [3], in the particular case where (i) the state space is $\mathbb{R}^{d}$ and (ii) the support of the jump measure is a finite set of $A$ different types of jumps. There, the approach is to study separately the jump-times of the reciprocal paths and their type distribution. In particular, the reciprocal characteristics come out analyzing the behaviour of the Poisson random vector $\mathbf{N} \in \mathbb{N}^{A}$ under shifts, where $\mathbf{N}$ describes the number of jumps of a path during the time interval $[0,1]$, classified by types.

In this paper we propose to characterize reciprocal classes of random walks in the following more general framework: the state space is a countable Abelian group $G$. See e.g. [17] for a review on random walks on groups.

Our tool is, by working directly at the level of the path space, to exhibit for each reciprocal class a family of integral equations which characterize it. The equations we obtain -see (14)- can be viewed as a twofold generalization of Mecke's formula, which characterizes Poisson random measures via transformations which consist in adding one point to the canonical point process. Indeed, first we iterate this procedure adding several jumps to the canonical process, and secondly we work under the constraint that the added paths are loops, that is they should have as initial and final value the identity element. However, our method is efficient only if one can assure that the set of loop paths is rich enough to allow to transform any given path of the random walk into any other one having the same initial and final value, only by adding and removing a finite number
of well chosen elementary loops. This assumption on the support $G_{\nu} \subseteq G$ of the jump measure $\nu$ is formalized through (H1) and (H2), see Section 2.2.

As an interesting byproduct of our integral characterization of the reciprocal classes of a $\nu$-random walk, we get the identification thanks equation (14) of the associated family of reciprocal characteristics (15). These quantities remain unchanged on the whole set of random walks having the same bridges, see Corollary 16.
The paper is organized as follows. In Section 1 we set up the necessary definitions and notations regarding random walks on groups, and provide an iterated formula on the path space satisfied by them. In Section 2, after having defined the concept of reciprocal classes, we state and prove our main result: the integral formula derived before on the path space is in fact a efficient way to characterize the whole reciprocal class of a random walk. In the last section we present several examples to illustrate the richness of the kind of processes we are dealing with.

## 1. Random walk on Abelian groups

### 1.1. The random walk as Poisson random measure

Let $(G,+)$ be a countable Abelian group with identity element $e$. We denote by $\mathbb{D}([0,1], G)$ the space of càdlàg paths for the topology induced by the discrete metric in $G$. Note that, because of the existence of left and right limits, paths in $\mathbb{D}([0,1], G)$ have finitely many jumps. $\mathbb{D}([0,1], G)$ is equipped with its canonical sigma-algebra $\mathcal{F}$ and its canonical filtration $\left(\mathcal{F}_{t}\right)_{t \in[0,1]}$.
For any $\nu$ non negative finite measure on $G$, we call $\nu$-random walk on $G$ a Markov probability on $\mathbb{D}([0,1], G)$ denoted by $\mathbb{P}_{\nu}$ whose infinitesimal generator $\mathscr{G}$ is given by

$$
(\mathscr{G} \phi)(g):=\sum_{g^{\prime} \in G} \nu\left(g^{\prime}\right)\left(\phi\left(g+g^{\prime}\right)-\phi(g)\right), g \in G
$$

for any $\phi$ bounded function. In the rest of the paper $G_{\nu} \subseteq G$ denotes the support of $\nu$, that is the set of allowed jumps of the $\nu$-random walk. The path space $\Omega \subset \mathbb{D}([0,1], G)$ is the set of paths with jumps in $G_{\nu}$.

Mecke proved in [11] an integral characterization of Poisson point processes on general spaces which we recall in Proposition 2 in a form adapted to our framework. In the spirit of Murr [13] and Privault ([14], Section 6.4.4) who studied real-valued processes with independent increments, we turn Mecke's formula into a characterization of random walks on $G$ in Proposition 6. Let us first introduce some notations.
For a measurable space $\mathcal{X}$ we denote by $\mathcal{P}(\mathcal{X})$ the set of probability measures on $\mathcal{X}$ and by $\mathcal{M}(\mathcal{X})$ the set of finite point measures, that is

$$
\mathcal{M}(\mathcal{X}):=\left\{\sum_{i=1}^{N} \delta_{x_{i}}: x_{i} \in \mathcal{X}, N \in \mathbb{N}\right\} .
$$

$\mathcal{B}^{+}(\mathcal{X})$ denotes the set of non-negative bounded measurable functions over $\mathcal{X}$. We will often choose for $\mathcal{X}$ the following product space $\Gamma$ of time-space elements:

$$
\Gamma:=[0,1] \times G \quad \ni \gamma=(t, g) .
$$

We identify trajectories in $\mathbb{D}([0,1], G)$ and point measures in $\mathcal{M}(\Gamma)$ via the following canonical bijective map $M$ :

$$
\begin{equation*}
X \mapsto M_{X}:=\sum_{0 \leq t \leq 1} \sum_{g \in G} \delta_{(t, g)} \mathbb{I}_{\left\{\Delta X_{t}=g\right\}} \tag{1}
\end{equation*}
$$

A useful observation is that the image measure of the $\nu$-random walk under $M$ is a Poisson random measure on $\Gamma$ with intensity the finite measure $d t \otimes \nu$.

### 1.2. An integral characterization

Mecke's original idea was to characterize any Poisson random measure by mean of an integral formula (see Satz 3.1 in [11]), via the change of measures which consists to add one (random) atom to the initial point measure, as in the right-hand side of Equation (3). Adapted to our context it reads as follows.

Proposition 2. For $\tilde{\mathbb{P}} \in \mathcal{P}(\mathcal{M}(\Gamma))$ the following assertions are equivalent:
i) $\tilde{\mathbb{P}}$ is the Poisson random measure with intensity measure $\rho=d t \otimes \nu$ on $\Gamma$.
ii) For all $\Phi \in \mathcal{B}^{+}(\mathcal{M}(\Gamma) \times \Gamma)$,

$$
\begin{equation*}
\iint_{\Gamma} \Phi(\mu, \gamma) \mu(d \gamma) \tilde{\mathbb{P}}(d \mu)=\iint_{\Gamma} \Phi\left(\mu+\delta_{\gamma}, \gamma\right) \rho(d \gamma) \tilde{\mathbb{P}}(d \mu) \tag{3}
\end{equation*}
$$

Remark that the left-hand side of (3) also reads $\int \sum_{\gamma \in \mu} \Phi(\mu, \gamma) \tilde{\mathbb{P}}(d \mu)$ where the notation $\gamma \in \mu$ means that the points $\gamma \in \Gamma$ build the support of the point measure $\mu$ : one integrates the function $\Phi$ under the Campbell measure associated with $\tilde{\mathbb{P}}$. Thus (3) determines the Campbell measure of a Poisson random measure as the shifted product measure of itself with its intensity.
Let us adapt this tool to $\mathbb{D}([0,1], G)$. First, for $\gamma=(t, g) \in \Gamma$, let us denote by $\chi^{\gamma}$ the corresponding simple step function $\chi^{\gamma}:=g \mathbb{\mathbb { I }}_{[t, 1]} \in \mathbb{D}([0,1], G)$. Then define the transformation $\Psi_{\gamma}$ on the path space which consists in adding one jump $g$ at time $t$.

Definition 4. For $\gamma=(t, g) \in \Gamma, \quad \Psi_{\gamma} X:=X+\chi^{\gamma}, \quad X \in \mathbb{D}([0,1], G)$.
Notice that, under any probability $\mathbb{P} \in \mathcal{P}(\mathbb{D}([0,1], G))$ satisfying $\mathbb{P}\left(X_{t}=X_{t^{-}}\right)=1$ for all $t \in[0,1]$, one has:

$$
\begin{equation*}
M_{\Psi_{\gamma} X}=M_{X}+\delta_{\gamma} \quad \mathbb{P}-a . s . \tag{5}
\end{equation*}
$$

Let us rewrite Proposition 2 in the language of random walks.

Proposition 6. For $\mathbb{P} \in \mathcal{P}(\mathbb{D}([0,1], G))$ the following assertions are equivalent:
i) $\mathbb{P}$ is a $\nu$-random walk on $G$.
ii) For all $F \in \mathcal{B}^{+}(\mathbb{D}([0,1], G) \times \Gamma)$,

$$
\begin{equation*}
\mathbb{P}\left(\int_{\Gamma} F(X, \gamma) M_{X}(d \gamma)\right)=\mathbb{P}\left(\int_{\Gamma} F\left(\Psi_{\gamma} X, \gamma\right) \rho(d \gamma)\right) \tag{7}
\end{equation*}
$$

where $M_{X}$ is defined through (1).
Proof. i) $\Rightarrow$ ii).
Since $\mathbb{P}_{\nu}$ is $\nu$-random walk, $M_{X}$ is a Poisson random measure with intensity $d t \otimes \nu$. Then Mecke's formula holds for $\tilde{\mathbb{P}}:=\mathbb{P} \circ M^{-1}$. Since $M$ is invertible and its inverse is measurable we can plug into (3) test functions $\Phi$ of the form $F(X, \gamma)$ and the conclusion follows.

$$
i i) \Rightarrow i)
$$

Let $\mathbb{P} \in \mathcal{P}(\mathbb{D}([0,1], G))$ satisfying (7). We define $\tilde{\mathbb{P}}:=\mathbb{P} \circ M^{-1} \in \mathcal{P}(\mathcal{M}(\Gamma))$. Then, by considering test functions of the form $\Phi=F\left(M_{X}, \gamma\right)$ and using the fact that $M_{\Psi_{\gamma} X}=$ $M_{X}+\delta_{\gamma} \mathbb{P} \otimes \rho-a . s .$, we deduce that $\tilde{\mathbb{P}}$ is a Poisson random measure with intensity $\rho=d t \otimes \nu$ by Proposition 2. Observing that

$$
X_{t}=\sum_{g \in G} g M_{X}([0, t] \times\{g\})
$$

the conclusion follows using (5).

### 1.3. An iterated formula

To prepare the characterization of bridges which we will present in the next section, we now consider a generalization of the formula (7) obtained by iteration. For this purpose, we define $n$-dimensional analogous of the objects appearing in (7), $n \geq 1$ fixed.

$$
\text { For } \bar{\gamma}=\left(\gamma_{1}, \ldots, \gamma_{n}\right) \in \Gamma^{n}, \Psi_{\bar{\gamma}}:=\Psi_{\gamma_{n}} \circ \ldots \circ \Psi_{\gamma_{1}}, \quad \chi^{\bar{\gamma}}=\sum_{i=1}^{n} \chi^{\gamma_{i}} .
$$

Then $\Psi_{\bar{\gamma}}$ is the transformation that adds to $X$ the $n$-step function $\chi^{\bar{\gamma}}$, which corresponds in the framework of point measures to add $n$ new atoms to the original measure. It remains to specify the set over which we will integrate the point measure $M_{X}^{\otimes n}$. Indeed, we remove from $\Gamma^{n}$ the set $\Delta_{n}$, union of all diagonals:

$$
\Delta_{n}:=\left\{\bar{\gamma} \in \Gamma^{n}: \exists i \neq j, \gamma_{i}=\gamma_{j}\right\} .
$$

With these definitions, we can state the iterated formula satisfied under the $\nu$-random walk.

Proposition 8. Let $\mathbb{P}_{\nu}$ be a $\nu$-random walk on $G$. Then, for any test function $F \in$ $\mathcal{B}^{+}\left(\mathbb{D}([0,1], G) \times \Gamma^{n}\right)$,

$$
\begin{equation*}
\mathbb{P}_{\nu}\left(\int_{\Gamma^{n} \backslash \Delta_{n}} F(X, \bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})\right)=\mathbb{P}_{\nu}\left(\int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right) \rho^{\otimes n}(d \bar{\gamma})\right) \tag{9}
\end{equation*}
$$

Remark 10. In general it is not true that $\int_{\Delta_{n}} F(X, \bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})=0$. Indeed, if $\gamma$ is an atom of $M_{X}$ then $\underbrace{(\gamma, \ldots, \gamma)}_{n \text { times }}$ is an atom of $M_{X}^{\otimes n}$ which belongs to $\Delta_{n}$.

Proof. We first prove a preliminary result.
Lemma 11. Define for $\gamma \in \Gamma, A_{\gamma}:=\left\{\bar{\gamma} \in \Gamma^{n}: \gamma_{i}=\gamma\right.$ for some $\left.1 \leq i \leq n\right\}$. Then

$$
M_{\Psi_{\gamma} X}^{\otimes n}(\cdot)=M_{X}^{\otimes n}(\cdot)+M_{\Psi_{\gamma} X}^{\otimes n}\left(A_{\gamma} \cap \cdot\right), \quad \mathbb{P}_{\nu} \otimes \rho a . e .
$$

Proof. We denote by $\mathcal{E}(X) \subset \Gamma$ the set of atoms of $M_{X}$ and by $\mathcal{E}^{n}(X)$ the set of atoms of $M_{X}^{\otimes n}$. Clearly $\mathcal{E}^{n}(X)=(\mathcal{E}(X))^{n} \subset \Gamma^{n}$. From (5),

$$
\mathcal{E}\left(\Psi_{\gamma} X\right)=\mathcal{E}(X) \cup\{\gamma\} \quad \mathbb{P}_{\nu} \otimes \rho \text { a.e. }
$$

so that the atoms of $M_{\Psi_{\gamma} X}^{\otimes n}$ are $\mathcal{E}^{n}\left(\Psi_{\gamma} X\right)=(\mathcal{E}(X) \cup\{\gamma\})^{n} \quad \mathbb{P}_{\nu} \otimes \rho$ a.e.. In this case $\mathcal{E}^{n}\left(\Psi_{\gamma} X\right) \backslash \mathcal{E}^{n}(X)=\mathcal{E}^{n}\left(\Psi_{\gamma} X\right) \cap A_{\gamma}$, which leads to the conclusion.

To prove the proposition by induction, we adopt the following notation: we decompose any element $\bar{\gamma}=\left(\gamma_{1}, \ldots, \gamma_{n+1}\right)$ of $\Gamma^{n+1}$ into $(\tilde{\gamma}, \gamma)$ where $\tilde{\gamma}$ is its projection on $\Gamma^{n}$, and $\gamma=\gamma_{n+1}$ is its last coordinate.
Proof of the statement for $n=1$ : It is exactly Proposition 6.
Let us now assume that the statement is true for $n$, that is (9) holds true for all test functions $f \in \mathcal{B}^{+}\left(\mathbb{D}([0,1], G) \times \Gamma^{n}\right)$. Let us now consider $v \in \mathcal{B}^{+}(\Gamma)$ and prove that (9) holds for any function of the form $F(X, \bar{\gamma})=f(X, \tilde{\gamma}) v(\gamma)$. The extension to a general $F \in \mathcal{B}^{+}\left(\mathbb{D}([0,1], G) \times \Gamma^{n+1}\right)$ will be then standard. We have:

$$
\mathbb{P}_{\nu}\left(\int_{\Gamma^{n+1}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right) \rho^{\otimes n+1}(d \bar{\gamma})\right)=\mathbb{P}_{\nu}\left(\int_{\Gamma} \int_{\Gamma^{n}} f\left(\Psi_{\gamma}\left(\Psi_{\tilde{\gamma}} X\right), \tilde{\gamma}\right) \rho^{\otimes n}(d \tilde{\gamma}) v(\gamma) \rho(d \gamma)\right)
$$

Exchanging the order of integration, and applying the inductive hypothesis to $f\left(\Psi_{\gamma} X, \tilde{\gamma}\right)$ allows to rewrite the right hand side of the last identity as:

$$
\int_{\Gamma} \mathbb{P}_{\nu}\left(\int_{\Gamma^{n} \backslash \Delta_{n}} f\left(\Psi_{\gamma} X, \tilde{\gamma}\right) M_{X}^{\otimes n}(d \tilde{\gamma})\right) v(\gamma) \rho(d \gamma)
$$

We can apply Lemma 11 to rewrite the former integral as:

$$
\int_{\Gamma} \mathbb{P}_{\nu}\left(\int_{\left(\Gamma^{n} \backslash \Delta_{n}\right) \backslash A_{\gamma}} f\left(\Psi_{\gamma} X, \tilde{\gamma}\right) M_{\Psi_{\gamma} X}^{\otimes n}(d \tilde{\gamma})\right) v(\gamma) \rho(d \gamma)
$$

We apply Proposition 6 to $(X, \gamma) \mapsto \int_{\left(\Gamma^{n} \backslash \Delta_{n}\right) \backslash A_{\gamma}} f(X, \tilde{\gamma}) M_{X}^{n}(d \tilde{\gamma}) v(\gamma)$ and we obtain

$$
\begin{aligned}
& \mathbb{P}_{\nu}\left(\int_{\Gamma}\left(\int_{\left(\Gamma^{n} \backslash \Delta_{n}\right) \backslash A_{\gamma}} f(X, \tilde{\gamma}) M_{X}^{\otimes n}(d \tilde{\gamma})\right) v(\gamma) M_{X}(d \gamma)\right) \\
& =\mathbb{P}_{\nu}\left(\int_{\left\{\bar{\gamma} ; \gamma \in \Gamma, \tilde{\gamma} \in\left(\Gamma^{n} \backslash \Delta_{n}\right) \backslash A_{\gamma}\right\}} F(X, \bar{\gamma}) M_{X}^{\otimes n+1}(d \bar{\gamma})\right)
\end{aligned}
$$

It is easy to see that $\left\{\bar{\gamma}=(\tilde{\gamma}, \gamma): \gamma \in \Gamma, \tilde{\gamma} \in\left(\Gamma^{n} \backslash \Delta_{n}\right) \backslash A_{\gamma}\right\}=\Gamma^{n+1} \backslash \Delta_{n+1}$ and the conclusion follows.

## 2. The reciprocal class and its characterization

### 2.1. Bridges of the random walk and their mixtures

First, consider the set of pairs $(x, y) \in G^{2}$ for which the bridge of the $\nu$-random walk is meaningful:

$$
S(\nu):=\left\{(x, y) \in G^{2}: \mathbb{P}_{\nu}\left(X_{1}=y \mid X_{0}=x\right)>0\right\} .
$$

Then, for $(x, y) \in S(\nu)$ the bridge $\mathbb{P}_{\nu}^{x y}$ between $x$ and $y$ is defined by

$$
\mathbb{P}_{\nu}^{x y}(\cdot):=\frac{\mathbb{P}_{\nu}\left(\cdot \cap\left\{X_{0}=x, X_{1}=y\right\}\right)}{\mathbb{P}_{\nu}\left(X_{0}=x, X_{1}=y\right)}
$$

We now define the reciprocal class associated with $\nu$ as the set of all possible mixtures of bridges of $\mathbb{P}_{\nu}$ :

$$
\operatorname{Rec}(\nu)=\left\{\mathbb{Q} \in \mathcal{P}(\Omega): \mathbb{Q}=\int_{S(\nu)} \mathbb{P}_{\nu}^{x y} \mathbb{Q}_{01}(d x d y)\right\}
$$

where $\mathbb{Q}_{01}$ denotes the joint marginal law of $\mathbb{Q}$ at times 0 and 1 . Let us note that for $(x, y)$ fixed in $S(\nu)$, the bridge $\mathbb{P}_{\nu}^{x y}$ itself belongs to the reciprocal class $\mathcal{R e c}(\nu)$. For a recent review on reciprocal processes, reciprocal classes and the analysis of basic examples we refer the reader e.g. to [10] or [15].

### 2.2. Loops and their skeletons

We call loop a path in $\mathbb{D}([0,1], G)$ that starts and ends at the identity element $e$. For each path $X \in \Omega$ we define its skeleton as the application $\varphi_{X}: G_{\nu} \rightarrow \mathbb{N}$ defined by:

$$
\varphi_{X}(g):=M_{X}([0,1] \times\{g\}) .
$$

Thus $\varphi_{X}(g)$ counts how many times the jump $g$ occurrs along the path X.
If $X$ is a loop, we observe that

$$
\sum_{g \in G_{\nu}} \varphi_{X}(g) g=e
$$

Therefore, as $X$ varies in the set of all possible loops, $\varphi_{X}$ varies in the set

$$
\begin{equation*}
\mathscr{L}^{+}:=\left\{\varphi \in \mathbb{N}^{G_{\nu}}: \sum_{g \in G_{\nu}} \varphi(g) g=e, \ell(\varphi)<+\infty\right\} \tag{12}
\end{equation*}
$$

where $\ell(\varphi):=\sum_{g \in G_{\nu}}|\varphi(g)|$ is the length of $\varphi$. Enlarging this set to the maps $\varphi$ with negatives values by considering

$$
\mathscr{L}:=\left\{\varphi \in \mathbb{Z}^{G_{\nu}}: \sum_{g \in G_{\nu}} \varphi(g) g=e, \ell(\varphi)<+\infty\right\}
$$

one recovers for $\mathscr{L}$ a lattice structure, which will be very useful. In particular $\mathscr{L}$ admits a basis $\mathscr{B}$. Suppose now that one can choose $\mathscr{B} \subset \mathscr{L}^{+}$, which is the case if the following assumption (H1) is satisfied:

$$
\begin{equation*}
\operatorname{Span}\left(\mathscr{L}^{+}\right)=\mathscr{L} \tag{H1}
\end{equation*}
$$

where $\operatorname{Span}\left(\mathscr{L}^{+}\right)$is, as usual, the set of all integer combinations of elements of $\mathscr{L}^{+}$. From now on, we fix such a basis $\mathscr{B}$.
To any $\varphi^{*} \in \mathscr{B}$ we can associate the - non empty - set of loops whose skeleton is $\varphi^{*}$ :

$$
\Omega_{e, \varphi^{*}}:=\left\{X \in \Omega: X_{0}=X_{1}=e \text { and } \varphi_{X}=\varphi^{*}\right\}
$$

These paths have exactly $\varphi^{*}(g)$ jumps of type $g$, for all $g \in G_{\nu}$.
Furthermore, we have to assume that each jump in $G_{\nu}$ belongs to (at least) the skeleton of one loop, that is, the following assumption holds:

$$
\begin{equation*}
\forall g \in G_{\nu} \text { there exists } \varphi \in \mathscr{L} \text { such that } \varphi(g)>0 \tag{H2}
\end{equation*}
$$

Note that w.l.o.g. we can assume that this skeleton $\varphi$ belongs indeed to the basis $\mathscr{B}$. As we shall see in Section 2.4, assumptions (H1) and (H2) allow a fruitful decomposition of the path space $\Omega$. Heuristically, one can transform one path into any other one having the same initial and final values, by subsequently adding and removing loops whose skeleton belongs to $\mathscr{B}$. However, let us first state our main result.

### 2.3. Main result: an integral characterization of the reciprocal class

In the next theorem we state that the identity (9) appeared in Proposition 8 is not only valid over the whole reciprocal class $\mathcal{R e c}(\nu)$ but indeed characterizes it, if one restricts the set of test functions $F$ to some well chosen subset.
For each skeleton $\varphi^{*}$ in the basis $\mathscr{B}$, consider the following set of test functions:

$$
\mathscr{H}_{\varphi^{*}}:=\left\{F \in \mathcal{B}^{+}\left(\mathbb{D}([0,1], G) \times \Gamma^{\ell\left(\varphi^{*}\right)}\right): F(X, \bar{\gamma}) \equiv \mathbb{1}_{\left\{\chi^{\bar{\gamma}} \in \Omega_{e, \varphi^{*}}\right\}} F(X, \bar{\gamma})\right\}
$$

Therefore, we will restrict our attention to perturbations of the sample paths consisting in adding a loop $\chi^{\gamma}$ whose skeleton is equal to $\varphi^{*}$. Now we are ready for stating and proving the main result.

Theorem 13. The probability measure $\mathbb{Q} \in \mathcal{P}(\Omega)$ belongs to the reciprocal class $\mathcal{R e c}(\nu)$ if and only if for any skeleton $\varphi^{*}$ in the basis $\mathscr{B}$ and for all test functions $F \in \mathscr{H}_{\varphi^{*}}$, we have:

$$
\begin{equation*}
\mathbb{Q}\left(\int_{\Gamma^{n} \backslash \Delta_{n}} F(X, \bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})\right)=\Phi_{\varphi^{*}}^{\nu} \mathbb{Q}\left(\int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right)(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right), \tag{14}
\end{equation*}
$$

where $n=\ell\left(\varphi^{*}\right), \Lambda:=\sum_{g \in G} \delta_{g}$ is the counting measure on $G$ and

$$
\begin{equation*}
\Phi_{\varphi^{*}}^{\nu}:=\prod_{g \in G_{\nu}} \nu(g)^{\varphi^{*}(g)} \in \mathbb{R}^{+} \tag{15}
\end{equation*}
$$

In particular, if (14) holds true under $\mathbb{Q}$ satisfying $\mathbb{Q}\left(X_{0}=x, X_{1}=y\right)=1$ for some $(x, y) \in S(\nu)$, then $\mathbb{Q}$ is nothing else but the bridge $\mathbb{P}_{\nu}^{x y}$.

The positive number $\Phi_{\varphi^{*}}^{\nu}$ is called the reciprocal characteristics associated to the jump measure $\nu$ and the skeleton $\varphi^{*}$.

Corollary 16. The reciprocal characteristics are invariants of the reciprocal class in the following sense. Let $\nu$ and $\mu$ two non negative finite measures on $G$ with the same support. The reciprocal classes $\mathcal{R e c}(\nu)$ and $\mathcal{R e c}(\mu)$ coincide if and only if their family of reciprocal characteristics coincide:

$$
\begin{equation*}
\Phi_{\varphi^{*}}^{\mu}=\Phi_{\varphi^{*}}^{\nu}, \quad \forall \varphi^{*} \in \mathscr{B} \tag{17}
\end{equation*}
$$

In that case the bridges of both $\nu$ - and $\mu$-random walk on $G$ coincide too.
Remark 18. There is a remarkable probabilistic interpretation of the reciprocal characteristics $\Phi_{\varphi^{*}}^{\nu}$ as the leading factor, in the short-time expansion, of the probability that the $\nu$-random walk follows a loop with skeleton $\varphi^{*}$. This is proven for Markov processes on graphs in the forthcoming paper [4].

### 2.4. Proof of the main theorem

Proof. ( $\Rightarrow$ )
We use, as main argument, the specific form of the density with respect to $\mathbb{P}_{\nu}$ of any probability measure in $\mathcal{R e c}(\nu)$ as it was proved in [3], Proposition 1.5:

$$
\mathbb{Q} \in \mathcal{R e c}(\nu) \Rightarrow \mathbb{Q} \ll \mathbb{P}_{\nu}, \quad \text { and } \quad \frac{d \mathbb{Q}}{d \mathbb{P}_{\nu}}=h\left(X_{0}, X_{1}\right) \text { for some } h: G \times G \rightarrow \mathbb{R}^{+}
$$

Take now any $F \in \mathscr{H}_{\varphi^{*}}$. Then, using Identity (9), the definition of $\mathscr{H}_{\varphi^{*}}$ and the fact that $\left(\Psi_{\bar{\gamma}} X\right)_{0}=X_{0},\left(\Psi_{\bar{\gamma}} X\right)_{1}=X_{1}+e=X_{1}$, one gets

$$
\begin{aligned}
\mathbb{Q}\left(\int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right)(d t \otimes \nu)^{\otimes n}(d \bar{\gamma})\right) & =\mathbb{Q}\left(\int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right) \prod_{g \in G_{\nu}} \nu(g)^{\varphi_{\chi^{\bar{\gamma}}}(g)}(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right) \\
& \stackrel{F \in \mathscr{H}_{\varphi^{*}}}{=} \Phi_{\varphi^{*}}^{\nu} \mathbb{Q}\left(\int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right)(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right) \\
& =\Phi_{\varphi^{*}}^{\nu} \mathbb{P}_{\nu}\left(h\left(X_{0}, X_{1}\right) \int_{\Gamma^{n}} F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right)(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right) \\
& =\Phi_{\varphi^{*}}^{\nu} \mathbb{P}_{\nu}\left(\int_{\Gamma^{n}} h\left(\left(\Psi_{\bar{\gamma}} X\right)_{0},\left(\Psi_{\bar{\gamma}} X\right)_{1}\right) F\left(\Psi_{\bar{\gamma}} X, \bar{\gamma}\right)(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right) \\
& =\mathbb{P}_{\nu}\left(\int_{\Gamma^{n} \backslash \Delta_{n}} h\left(X_{0}, X_{1}\right) F(X, \bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})\right) \\
& =\mathbb{Q}\left(\int_{\Gamma^{n} \backslash \Delta_{n}} F(X, \bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})\right)
\end{aligned}
$$

which completes the proof of the first implication.

$$
(\Leftarrow)
$$

The converse implication is more sophisticated and needs several steps.
Let us introduce the set of paths which correspond to the support of ey-bridges, $y \in G$ :

$$
\Omega_{y}:=\left\{X \in \Omega: X_{0}=e, X_{1}=y\right\}
$$

Now we partition $\Omega_{y}$ according to the skeleton of its elements:

$$
\begin{gathered}
\Omega_{y}=\bigcup_{\varphi \in \mathscr{L}_{y}^{+}} \Omega_{y, \varphi}, \quad \Omega_{y, \varphi}:=\Omega_{y} \cap\left\{X \in \Omega: \varphi_{X}=\varphi\right\} \\
\text { where } \mathscr{L}_{y}^{+}=\left\{\varphi \in \mathbb{N}^{G_{\nu}}: \sum_{g \in G_{\nu}} g \varphi(g)=y, \ell(\varphi)<+\infty\right\}
\end{gathered}
$$

In order to discretize the time, we introduce a mesh $h \in \mathbb{N}^{*}$ and partition $\Omega_{y, \varphi}$ by specifying the number of different jumps occurred in each $h$-dyadic interval. That is, we consider functions $\theta:\left\{0, . ., 2^{h}-1\right\} \times G_{\nu} \longrightarrow \mathbb{N}$ and we look for paths which have $\theta(k, g)$ jumps of type $g$ during the time interval $I_{k}^{h}:=\left(2^{-h} k, 2^{-h}(k+1)\right.$ ], for each $k$ and each $g \in G_{\nu}$. For each skeleton $\varphi$ we define the set

$$
\Theta_{\varphi}^{h}:=\left\{\theta:\left\{0, . ., 2^{h}-1\right\} \times G_{\nu} \longrightarrow \mathbb{N}, \sum_{0 \leq k \leq 2^{h}-1} \theta(k, g)=\varphi(g), \quad \forall g \in G_{\nu}\right\}
$$

of all possible $h$-dyadic time repartition of the jumps, compatible with the skeleton $\varphi$. We thus obtain $\Omega_{y, \varphi}=\bigcup_{\theta \in \Theta_{\varphi}^{h}} \Omega_{y, \varphi}^{h, \theta}$ where

$$
\begin{aligned}
& \qquad \Omega_{y, \varphi}^{h, \theta}:=\left\{X \in \Omega_{y}: M_{X}\left(I_{k}^{h} \times\{g\}\right)=\theta(k, g), 0 \leq k<2^{h}, g \in G_{\nu}\right\} \\
& \text { imsart-bj ver. 2014/01/08 file: PreprintUP_Conforti_Roelly.tex date: January 6, } 2015
\end{aligned}
$$

Consider the set

$$
\begin{equation*}
V:=\left\{v=(\varphi, \theta) \text { with } \varphi \in \mathscr{L}_{y}^{+}, \theta \in \Theta_{\varphi}^{h}\right\} \tag{19}
\end{equation*}
$$

of pairs of skeletons connecting $e$ to $y$ and $h$-dyadic time repartition of their jumps. Elements of this set are discrete versions of paths of $\Omega$ : the spatial structure of the path is given by the skeleton $\varphi$, and the time structure is approximated by $\theta$. One equips $V$ with the following $l^{1}$-metric:

$$
d(v, \tilde{v}):=\sum_{(k, g) \in\left\{0, \cdots, 2^{h}-1\right\} \times G_{\nu}}|\theta-\tilde{\theta}|(k, g) \in \mathbb{N}, \quad v=(\varphi, \theta), \tilde{v}=(\tilde{\varphi}, \tilde{\theta}) \in V
$$

Take now two paths $X, X^{\prime} \in \Omega_{y}$ and their trace $v, v^{\prime}$ on V . Our aim is to find a way to transform $X$ into $X^{\prime}$ (resp. $v$ into $v^{\prime}$ ) by adding or removing a finite number of loops whose skeletons belong to the basis $\mathscr{B}$. Let us introduce the following relation:

$$
v_{1}=\left(\varphi_{1}, \theta_{1}\right) \hookrightarrow v_{2}=\left(\varphi_{2}, \theta_{2}\right) \text { if } \varphi_{2} \in \varphi_{1}+\mathscr{B} \text { and } \theta_{2}-\theta_{1} \in \Theta_{\varphi_{2}-\varphi_{1}}^{h}
$$

We shall now use assumptions (H1) and (H2).
Lemma 20. For each $v$ and $\tilde{v} \neq v \in V$ on can construct a connecting finite sequence $v_{1}, \cdots, v_{N}=\tilde{v}_{\tilde{N}}, \tilde{v}_{\tilde{N}-1}, \cdots, \tilde{v}_{1}$ in $V$ such that

$$
v \hookrightarrow v_{1} \hookrightarrow \cdots \hookrightarrow v_{N}=\tilde{v}_{\tilde{N}} \hookleftarrow \tilde{v}_{\tilde{N}-1} \cdots \hookleftarrow \tilde{v}_{1} \hookleftarrow \tilde{v}
$$

Proof. We distinguish two cases:
Case i) The skeletons $\varphi$ and $\tilde{\varphi}$ coincide.
In this case, it is sufficient to show that we can construct $v_{1}$ and $\tilde{v}_{1}$ in $V$ such that $v \hookrightarrow v_{1}, \tilde{v} \hookrightarrow \tilde{v}_{1}, \varphi_{1}=\tilde{\varphi}_{1}$ and $d\left(v_{1}, \tilde{v}_{1}\right) \leq d(v, \tilde{v})-1$. The conclusion would then follow by iterating this procedure until $d\left(v_{K}, \tilde{v}_{K}\right)=0$, i.e. $v_{K}=\tilde{v}_{K}$. At this point, we have constructed a chain from $v$ to $v_{K}$, and another one from $\tilde{v}$ to $\tilde{v}_{K}$. Joining them, we obtain a chain from $v$ to $\tilde{v}$ and the conclusion follows.
Therefore, let us indicate how to construct $v_{1}$ and $\tilde{v}_{1}$. Since $\theta \neq \tilde{\theta}$ but $\varphi=\tilde{\varphi}$ there exists a jump $g \in G_{\nu}$ and two time intervals $I_{k}^{h}$ and $I_{l}^{h}$ such that $\theta(k, g) \geq \tilde{\theta}(k, g)+1$ and $\theta(l, g) \leq \tilde{\theta}(l, g)-1$. Moreover, thanks to (H2) there exists at least one skeleton $\varphi^{*}$ in the basis $\mathscr{B}$ containing the jump $g: \varphi^{*}(g)>0$. Consider now any time repartition $\theta_{1} \in \Theta_{\varphi^{*}}^{h}$ such that $\theta_{1}(l, g) \geq 1$. We then construct $\tilde{\theta}_{1}$ as follows:

$$
\tilde{\theta}_{1}=\theta_{1}+\mathbb{1}_{\{(k, g)\}}-\mathbb{1}_{\{(l, g)\}} .
$$

It is simple to check that $v_{1}:=\left(\varphi+\varphi^{*}, \theta+\theta_{1}\right), \tilde{v}_{1}:=\left(\varphi+\varphi^{*}, \tilde{\theta}+\tilde{\theta}_{1}\right)$ fulfill the desired requirements. By construction, $v \hookrightarrow v_{1}, \tilde{v} \hookrightarrow \tilde{v}_{1}$ and $v_{1}, \tilde{v}_{1}$ have the same skeleton. Moreover

$$
\left|\theta+\theta_{1}-\left(\tilde{\theta}+\tilde{\theta}_{1}\right)\right|=|\theta-\tilde{\theta}|-\mathbb{1}_{\{(k, g),(l, g)\}}
$$

so that $d\left(v_{1}, \tilde{v}_{1}\right)=d(v, \tilde{v})-2$.

Case ii) The skeletons $\varphi$ and $\tilde{\varphi}$ differ.
We first observe that, if $\varphi, \tilde{\varphi} \in \mathscr{L}_{y}^{+}$thus $\varphi-\tilde{\varphi} \in \mathscr{L}$. Since $\mathscr{B}$ is a basis of the lattice $\mathscr{L}($ see $(\mathrm{H} 1))$, there exist $\left(\varphi_{j}^{*}\right)_{j=1}^{K},\left(\tilde{\varphi}_{i}^{*}\right)_{i=1}^{\tilde{K}} \subseteq \mathscr{B}$ such that

$$
\varphi+\sum_{j=1}^{K} \varphi_{j}^{*}=\tilde{\varphi}+\sum_{i=1}^{\tilde{K}} \tilde{\varphi}_{i}^{*}
$$

Let us now choose for all $j$ and $i$ a time repartition $\theta_{j} \in \Theta_{\tilde{\varphi}_{j}^{*}}^{h}$ and $\tilde{\theta}_{i} \in \Theta_{\tilde{\varphi}_{i}}^{h}$. It is straightforward to verify that, if we define

$$
v_{0}=v, v_{j}:=\left(\varphi+\sum_{j^{\prime}=1}^{j} \varphi_{j^{\prime}}^{*}, \theta+\sum_{j^{\prime}=1}^{j} \theta_{j^{\prime}}\right), \tilde{v}_{0}=\tilde{v}, \tilde{v}_{i}:=\left(\tilde{\varphi}+\sum_{i^{\prime}=1}^{i} \tilde{\varphi}_{i^{\prime}}^{*}, \tilde{\theta}^{*}+\sum_{i^{\prime}=1}^{i} \tilde{\theta}_{i^{\prime}}\right)
$$

then $\left(v_{j}\right)_{j=0}^{K},\left(\tilde{v}_{i}\right)_{i=0}^{\tilde{K}}$ are two sequences connecting $v$ to $v_{K}$ and $\tilde{v}$ to $\tilde{v}_{\tilde{K}}$. By construction $v_{K}, \tilde{v}_{\tilde{K}}$ have the same skeleton and one can use case i) again.


Figure 1. In this picture we illustrate by an example the proof of Lemma 20. Take $G=(\mathbb{Z},+)$, and $G_{\nu}=\{-1,1,2\}$, situation which is treated in Section 3.1.1. $\mathscr{B}=\left\{\varphi_{1}, \varphi_{2}\right\}$, where $\varphi_{1}:=\mathbb{I}_{1}+\mathbb{I}_{-1}$ and $\varphi_{2}:=\mathbb{I}_{2}+2 \mathbb{I}_{-1}$, is a basis fulfilling H1) and H2). The picture shows how to transform the path a) in the path f ) by mean of addition and cancellation of loops whose skeleton belongs to $\mathscr{B}$. All loops that are either added or removed are denoted by red dashed lines, which correspond to their jumps. At first, following case ii), we have to modify the loop a) to match its skeleton $(2,2,0)$ with that of f$),(3,1,1)$. Therefore in b) we remove a loop with skeleton $\varphi_{1}$, then in c) add back a loop with skeleton $\varphi_{2}$. The skeleton is now the desired one. Now we follow case i): we shift one jump of height -1 and one of height 1 further right. Since those two jumps form a loop with skeleton $\varphi_{1}$ we simply delete them in d) and add a new loop with the same skeleton, but now with the desired jump times in e).

In the next lemma we compare the probability of the paths in $\Omega_{y, \varphi}^{h, \theta}$ and those obtained by adding a loop with skeleton $\varphi^{*} \in \mathscr{B}$, under $\mathbb{Q}$ and under $\mathbb{P}_{\nu}$.

Lemma 21. Let $y \in G, h \in \mathbb{N}^{*}, \varphi \in \mathscr{L}_{y}^{+}, \theta \in \Theta_{\varphi}^{h}$ be fixed. Suppose (14) holds under $\mathbb{Q}$. Then, for any $\varphi^{*} \in \mathscr{B}$ and $\theta^{*} \in \Theta_{\varphi^{*}}^{h}$,

$$
\begin{equation*}
\frac{\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta+\theta^{*}}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}\right)}=\frac{\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right)} \tag{22}
\end{equation*}
$$

Proof. Take an arbitrary ordering of the support of $\theta^{*}:\left(k_{1}, g_{1}\right), \ldots,\left(k_{N}, g_{N}\right)$. To simplify the notation, we write $\theta_{j}\left(\right.$ resp. $\left.\theta_{j}^{*}\right)$ for $\theta\left(k_{j}, g_{j}\right)$ (resp $\left.\theta^{*}\left(k_{j}, g_{j}\right)\right)$. Consider the test function $F(X, \bar{\gamma})=f(X) v(\bar{\gamma})$, where

$$
f=\mathbb{1}_{\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}}, \text { and } v(\bar{\gamma})=\mathbb{1}_{\Omega_{e, \varphi^{*}}^{h, \theta^{*}}}\left(\chi^{\bar{\gamma}}\right) .
$$

It is straightforward that

$$
f \circ \Psi_{\bar{\gamma}}(X) v(\bar{\gamma})=\mathbb{1}_{\Omega_{y, \varphi}^{h, \theta}}(X) v(\bar{\gamma}) \quad \mathbb{Q} \otimes \rho^{n} \text { a.e.. }
$$

Therefore, since $F \in \mathscr{H}_{\varphi^{*}}$, (14) holds and its right hand side rewrites as

$$
\Phi_{\varphi^{*}}^{\nu}\left(\int_{\Gamma^{n}} v(\bar{\gamma})(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma})\right) \mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)
$$

Concerning the left hand side, let us first rewrite it as

$$
\mathbb{Q}\left(f(X) \int_{\Gamma^{n} \backslash \Delta_{n}} v(\bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})\right) .
$$

Our aim is to show by a direct computation that the (discrete) stochastic integral $\int_{\Gamma^{n} \backslash \Delta_{n}} v(\bar{\gamma}) d M_{X}^{\otimes n}(d \bar{\gamma})$ is actually constant for that choice of $v$ if $X \in \Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}$.
First, we observe that an atom $\bar{\gamma} \in \Gamma^{n} \backslash \Delta_{n}$ of $M_{X}^{\otimes n}$ contributes (with the value 1) to the integral if and only if $\chi^{\bar{\gamma}} \in \Omega_{e, \varphi^{*}}^{h, \theta^{*}}$, that is if

$$
\begin{equation*}
\sharp\left\{i: \gamma_{i} \in I_{k_{j}}^{h} \times\left\{g_{j}\right\}\right\}=\theta_{j}^{*}, \quad 1 \leq j \leq N . \tag{23}
\end{equation*}
$$

We then need to count the atoms of $M_{X}^{\otimes n}$ satisfying (23). This is equivalent to count all ordered lists of $n=\ell\left(\varphi^{*}\right)$ atoms of $M_{X}$ verifying that

1) the list contains no repetitions
2) for all $1 \leq j \leq N$, the number of elements in the list which belong to $I_{k_{j}}^{h} \times\left\{g_{j}\right\}$ is $\theta_{j}^{*}$. Therefore, for each $j$, we first choose a subset of cardinality $\theta_{j}^{*}$ among $\theta_{j}+\theta_{j}^{*}$ elements (recall that $X \in \Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}$ ). To do that we have ( $\binom{\theta_{j}+\theta_{j}^{*}}{\theta_{j}^{*}}$ choices. Then we decide how to sort the list, and for this, there are $n$ ! possibilities.
Therefore

$$
\mathbb{I}_{\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}}(X) \int_{\Gamma^{n} \backslash \Delta_{n}} v(\bar{\gamma}) M_{X}^{\otimes n}(d \bar{\gamma})=\mathbb{1}_{\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}}(X) n!\prod_{j=1}^{N}\binom{\theta_{j}+\theta_{j}^{*}}{\theta_{j}^{*}}
$$

and (14) rewrites as

$$
\begin{equation*}
\Phi_{\varphi^{*}}^{\nu} \int_{\Gamma^{n}} v(\bar{\gamma})(d t \otimes \Lambda)^{\otimes n}(d \bar{\gamma}) \mathbb{Q}\left(\Omega_{y, \varphi}\right)=n!\prod_{j=1}^{N}\binom{\theta_{j}+\theta_{j}^{*}}{\theta_{j}^{*}} \mathbb{Q}\left(\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}\right) \tag{24}
\end{equation*}
$$

Since equation (14) holds under $\mathbb{P}_{\nu}$, equation (24) holds under $\mathbb{P}_{\nu}$ as well. Since $\mathbb{P}_{\nu}$ gives positive probability to both events $\Omega_{y, \varphi}$ and $\Omega_{y, \varphi+\varphi^{*}}^{h, \theta+\theta^{*}}$, the identity (22) follows.

Remark that, with the notation of the above lemma, if we define $v:=(\varphi, \theta)$ and $w:=\left(\varphi+\varphi^{*}, \theta+\theta^{*}\right)$, then $v \hookrightarrow w$.

Lemma 20 allows us to extend the conclusion of Lemma 21 to the whole set of skeletons, as we will prove now.

Lemma 25. Let $y \in G, h \in \mathbb{N}^{*}, \varphi, \tilde{\varphi} \in \mathscr{L}_{y}^{+}, \theta \in \Theta_{\varphi}^{h}, \tilde{\theta} \in \Theta_{\tilde{\varphi}}^{h}$ be fixed. Suppose (14) holds under $\mathbb{Q}$. Then,

$$
\begin{equation*}
\frac{\mathbb{Q}\left(\Omega_{y, \tilde{\tilde{\theta}}}^{h, \tilde{\theta}}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \tilde{\varphi}}^{h, \tilde{\theta}}\right)}=\frac{\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right)} \tag{26}
\end{equation*}
$$

Proof. We observe that $v=(\varphi, \theta)$ and $\tilde{v}=(\tilde{\varphi}, \tilde{\theta})$ are elements of $V$. As proved above, there exists a connecting sequence $\left(v_{i}\right)_{i=0}^{K}:=\left(\varphi_{i}, \theta_{i}\right)_{i=0}^{K}$, with $v_{0}=v, v_{K}=\tilde{v}$, linking $v$ to $\tilde{v}$, and such that either $v_{i} \hookleftarrow v_{i+1}$ or $v_{i} \hookrightarrow v_{i+1}$. This entitles us to apply recursively Lemma 21 to any pair $v_{i}, v_{i+1}$ and obtain

$$
\frac{\mathbb{Q}\left(\Omega_{y, \varphi_{i+1}}^{h, \theta_{i+1}}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi_{i+1}}^{h, \theta_{i+1}}\right)}=\frac{\mathbb{Q}\left(\Omega_{y, \varphi_{i}}^{h, \theta_{i}}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi_{i}}^{h,,_{i}}\right)}=\ldots=\frac{\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)}{\mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right)}
$$

The conclusion follows with $i=N-1$.

We can now complete the proof of the converse implication of the main theorem.
Fix $x, y \in G$ with $\mathbb{Q}\left(X_{0}=x, X_{1}=y\right)>0$. W.l.o.g. we assume that $x=e$. Thanks to Lemma 25 we know that for any mesh $h$, there exists a positive constant $c_{h}$ such that

$$
\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)=c_{h} \mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right), \quad \forall \varphi \in \mathscr{L}_{y}^{+}, \theta \in \Theta_{\varphi}^{h}
$$

Now we show that the proportionality constant does not depend on the scale of the time discretisation: $c_{h}=c_{h+1}$. To this aim, let us observe that

$$
\mathbb{Q}\left(\Omega_{y}\right)=\sum_{(\varphi, \theta) \in V} \mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)=\sum_{(\varphi, \theta) \in V} c_{h} \mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right)=c_{h} \mathbb{P}_{\nu}\left(\Omega_{y}\right)
$$

In the same way one gets $\mathbb{Q}\left(\Omega_{y}\right)=c_{h+1} \mathbb{P}_{\nu}\left(\Omega_{y}\right)$ which implies that $c_{h}=c_{h+1}$. Therefore, there exists a constant $c>0$ such that

$$
\mathbb{Q}\left(\Omega_{y, \varphi}^{h, \theta}\right)=c \mathbb{P}_{\nu}\left(\Omega_{y, \varphi}^{h, \theta}\right), \quad \forall h \in \mathbb{N}^{*}, \varphi \in \mathscr{L}_{y}^{+}, \theta \in \Theta_{\varphi}^{h}
$$

By standard approximation arguments this implies the equality between $\mathbb{Q}$ and $c \mathbb{P}_{\nu}$ on $\Omega_{y} \cap \mathcal{F}$ which then implies $\mathbb{Q}^{e y}=\mathbb{P}_{\nu}^{e y}$. The conclusion follows.

Remark 27. Consider the identities (14) for $G=\mathbb{R}^{d}$ and compute them for particular test functions $F$ which only depend on the skeleton of the paths. These equations, indexed by the skeletons in $\mathscr{B}$, then characterize the (marginal) distribution of the random vector defined as the number of jumps of any type occurred during the time interval $[0,1]$, as it was done in [3]. Note that for the unconstrained random walk the distribution of this random vector is a multivariate Poisson law, see e.g. [3] Section 2.2.1.

## 3. Examples

In this section, we present several examples of random walks defined on finite or infinite Abelian groups $G$.
For each example, we verify if assumptions (H1) and (H2) are satisfied by computing a basis $\mathscr{B}$ of skeleton of loops. We give explicitly the associated reciprocal characteristics (15). In some cases, we also write down the integral formula (14), highlighting how it is influenced by the geometrical properties of the underlying group $G$.
Finally, for a fixed random walk $\mathbb{P}_{\nu}$ on $G$, we address the question of finding all random walks $\mathbb{P}_{\mu}$ which have the same bridges than $\mathbb{P}_{\nu}$, that is, using Corollary 16, we solve equation (17) and identify the set of probability measures:

$$
\mathcal{R e c}(\nu) \cap\left\{\mathbb{P}_{\mu}: \mu \text { finite measure on } G_{\nu}\right\} .
$$

We will see that, in some cases, this set reduces to the singleton $\mathbb{P}_{\nu}$ and in other cases, this set is non trivial.

### 3.1. The group $G=\mathbb{Z}$ is infinite

3.1.1. The finite support $G_{\nu}$ of the jump measure $\nu$ contains $\{-1,1\}$.

For any $z \in G_{\nu} \backslash\{1\}$ we define on $G_{\nu}$ the non negative map $\varphi_{z}$ as follows:

$$
\varphi_{z}=\mathbb{1}_{z}+|z| \mathbb{I}_{-\operatorname{sgn}(z)}
$$

It corresponds to the skeleton of paths with one jump of type $z$ and $|z|$ jumps of type $-\operatorname{sgn}(z)$. As candidate for the lattice basis of $\mathscr{L}$, we propose

$$
\mathscr{B}:=\left\{\varphi_{z}\right\}_{z \in G_{\nu} \backslash\{1\}} .
$$

Assumption (H2) is trivially satisfied and it is clear that the elements of $\mathscr{B}$ are linearly independent. Therefore we only need to check if $\mathscr{B}$ spans $\mathscr{L}$, that is, if for each $\phi \in \mathscr{L}$, there exist integer coefficients $\alpha_{z} \in \mathbb{Z}$ such that

$$
\begin{equation*}
\forall \bar{z} \in G_{\nu}, \quad \phi(\bar{z})=\sum_{\substack{z \in G_{\nu} \\ z \neq 1}} \alpha_{z} \varphi_{z}(\bar{z}) . \tag{28}
\end{equation*}
$$

We now verify that the following choice is the right one:

$$
\text { For } z \in G_{\nu} \backslash\{-1,+1\} \alpha_{z}=\phi(z) \text { and } \alpha_{-1}=\phi(-1)-\sum_{\substack{z \in G_{\nu} \\ z>1}} z \phi(z)
$$

- $\bar{z} \notin\{-1,+1\}$. Since $\varphi_{\bar{z}}$ is the only element of $\mathscr{B}$ whose support contains $\bar{z}$, we have

$$
\phi(\bar{z})=\alpha_{\bar{z}} \varphi_{\bar{z}}(\bar{z})=\sum_{\substack{z \in G_{\nu} \\ z \neq 1}} \alpha_{z} \varphi_{z}(\bar{z})
$$

- $\bar{z}=-1$. Notice that -1 belongs to the support of any $\varphi_{z}$, as soon as $z>1$. Therefore
$\phi(-1)=\sum_{\substack{z \in G_{\nu} \\ z>1}} \phi(z) z+\alpha_{-1}=\sum_{\substack{z \in G_{\nu} \\ z>1}} \alpha_{z} \varphi_{z}(-1)+\alpha_{-1} \varphi_{-1}(-1)=\left(\sum_{\substack{z \in G_{\nu} \\ z \neq 1}} \alpha_{z} \varphi_{z}\right)(-1)$.
- $\bar{z}=1$. Notice that +1 belongs to the support of any $\varphi_{z}$, as soon as $z \leq-1$. Recall that $\phi \in \mathscr{L}$. Therefore

$$
\begin{aligned}
\phi(1)=-\sum_{\substack{z \in G_{\nu} \\
z \neq 1}} \phi(z) z & =\sum_{\substack{z \in G_{\nu} \\
z<1}}-\phi(z) z+\phi(-1) \\
& =\sum_{\substack{z \in G_{\nu} \\
z \leq-1}} \alpha_{z} \varphi_{z}(1)=\left(\sum_{\substack{z \in G_{\nu} \\
z \neq 1}} \alpha_{z} \varphi_{z}\right)(1) .
\end{aligned}
$$

Let us now compute the reciprocal characteristics associated to each skeleton in $\mathscr{B}$ :

$$
\Phi_{\varphi_{z}}^{\nu}=\nu(-\operatorname{sgn}(z))^{|z|} \nu(z), \quad z \in G_{\nu} \backslash\{1\}
$$

Finally, thanks to Corollary 16, we obtain

$$
\begin{aligned}
\mathbb{P}_{\mu} \in \operatorname{Rec}(\nu) & \Leftrightarrow \forall z \in G_{\nu} \backslash\{1\}, \quad \mu(-\operatorname{sgn}(z))^{|z|} \mu(z)=\nu(-\operatorname{sgn}(z))^{|z|} \nu(z) \\
& \Leftrightarrow \exists \alpha>0 \text { such that } \frac{d \mu}{d \nu}(z)=\alpha^{z}
\end{aligned}
$$

Example 29. Simple random walks: $G_{\nu}=\{-1,1\}$.
Due to the above computations, the basis $\mathscr{B}$ of the lattice $\mathscr{L}$ reduces to the singleton $\left\{\varphi_{-1}\right\}$ and the unique reciprocal characteristics is given by

$$
\Phi_{\varphi_{-1}}^{\nu}=\nu(-1) \nu(1)
$$

Therefore the only loops which appear in the integral characterization (14) have length $n=\ell\left(\varphi_{-1}\right)=2$.
Test functions of the form

$$
F\left(X,\left(\gamma_{1}, \gamma_{2}\right)\right)=f(X) \mathbb{I}_{\left\{g_{1}=1, g_{2}=-1\right\}} h\left(t_{1}, t_{2}\right)
$$

belong to $\mathscr{H}_{\varphi_{-1}}$. Such functions are supported by pairs $\left(\gamma_{1}, \gamma_{2}\right)$ such that the transformation $\Psi_{\gamma_{1}, \gamma_{2}}$ adds to any path a jump +1 at time $t_{1}$ and a jump-1 at time $t_{2}$. The identity (14) now reads as:

$$
\mathbb{Q}\left(f(X) \sum_{\substack{\left(t_{1}, t_{2}\right): \Delta X_{t_{1}}=1 \\ \Delta X_{t_{2}}=-1}} h\left(t_{1}, t_{2}\right)\right)=\nu(-1) \nu(1) \int_{[0,1]^{2}} \mathbb{Q}\left(f\left(\Psi_{\gamma_{1}, \gamma_{2}} X\right)\right) h\left(t_{1}, t_{2}\right) d t_{1} d t_{2}
$$

As in Remark 27, if we consider test functions $f$ which only depend on the skeletons of the paths, $f(X)=v\left(\varphi_{X}\right)$, we obtain that the distribution $\chi\left(d n_{-1}, d n_{1}\right) \in \mathcal{P}\left(\mathbb{N}^{2}\right)$ of the
number $n_{-1}$ (resp. $n_{1}$ ) of negative (resp. positive) jumps is characterized by the system of equations: for all $v \in \mathcal{B}^{+}\left(\mathbb{N}^{2}\right)$,

$$
\begin{aligned}
\int v\left(n_{-1}, n_{1}\right) n_{-1} n_{1} \chi\left(d n_{-1}, d n_{1}\right) & =\nu(-1) \nu(1) \int v\left(n_{-1}+1, n_{1}+1\right) \chi\left(d n_{-1}, d n_{1}\right) \\
\chi\left(n_{1}=n_{-1}\right) & =1
\end{aligned}
$$

This result coincides with [3], Example 2.18.
3.1.2. $G_{\nu}=\{1,2\}$.

In that case, since -1 does not belong to the support of the jump measure, it leads to a case where (H2) is not satisfied. It is straightforward to prove that the lattice $\mathscr{L}$ is one-dimensional and is equal to $\left\{z \varphi^{*}, z \in \mathbb{Z}\right\}$ where

$$
\varphi^{*}(1)=2, \quad \varphi^{*}(2)=-1
$$

Clearly $\mathscr{L}$ does not admit a non negative basis.

## 3.2. $G$ is the cyclic group $\mathbb{Z} / N \mathbb{Z}$

We now consider the finite cyclic group $G:=\mathbb{Z} / N \mathbb{Z}=:\{\mathbf{0}, \mathbf{1}, \mathbf{2}, \cdots, \mathbf{N}-\mathbf{1}\}$.

### 3.2.1. The support $G_{\nu}$ of the jump measure reduces to $\{-\mathbf{1}, \mathbf{1}\}$.

This case corresponds to nearest neighbour random walks. The non negative basis $\mathscr{B}:=$ $\left\{\varphi_{N-1}, \varphi^{*}\right\}$ where

$$
\varphi_{N-1}=\mathbb{I}_{1}+\mathbb{I}_{\mathbf{N - 1}}=\mathbb{1}_{1}+\mathbb{I}_{-\mathbf{1}} \text { and } \varphi^{*}=N \mathbb{1}_{1}
$$

is suitable. The associated reciprocal characteristics are

$$
\Phi_{\varphi_{N-1}}^{\nu}=\nu(\mathbf{1}) \nu(-\mathbf{1}) \text { and } \Phi_{\varphi^{*}}^{\nu}=\nu(\mathbf{1})^{N} .
$$

The existence of the second invariant $\Phi_{\varphi^{*}}^{\nu}$ corresponding to the loop around the cycle $\{\mathbf{0}, \mathbf{1}, \mathbf{2}, \cdots, \mathbf{N}-\mathbf{1}\}$ implies that $\mathbb{P}_{\nu}$ is the unique nearest neighbour random walk of the reciprocal class $\mathcal{R e c}(\nu)$. This differs from the nearest neighbour random walk on $\mathbb{Z}$, treated in Example 29. We proved there that any random walk $\mathbb{P}_{\mu}$, with $\mu$ satisfying $\mu(-1) \mu(1)=\nu(-1) \nu(1)$, induces the same reciprocal class.

The distribution $\chi$ of the random vector $\left(n_{-\mathbf{1}}, n_{\mathbf{1}}\right)$ under the $\mathbf{0 0}$-bridge is given by the following system of integral equations, satisfied for any test function $v$ on $\mathbb{N}^{2}$ :

$$
\begin{aligned}
\int v\left(n_{-1}, n_{\mathbf{1}}\right) n_{-\mathbf{1}} n_{\mathbf{1}} \chi\left(d n_{-\mathbf{1}}, d n_{\mathbf{1}}\right) & =\nu(-\mathbf{1}) \nu(\mathbf{1}) \int v\left(n_{-1}+1, n_{\mathbf{1}}+1\right) \chi\left(d n_{-\mathbf{1}}, d n_{\mathbf{1}}\right), \\
\int v\left(n_{-1}, n_{\mathbf{1}}\right) n_{\mathbf{1}} \cdots\left(n_{\mathbf{1}}-(N-1)\right) \chi\left(d n_{-\mathbf{1}}, d n_{\mathbf{1}}\right) & =\nu(\mathbf{1})^{N} \int v\left(n_{-\mathbf{1}}, n_{\mathbf{1}}+N\right) \chi\left(d n_{-\mathbf{1}}, d n_{\mathbf{1}}\right), \\
\chi\left(n_{\mathbf{1}}-n_{-\mathbf{1}} \in N \mathbb{Z}\right) & =1 .
\end{aligned}
$$

### 3.2.2. The support $G_{\nu}$ of the jump measure covers $\mathbb{Z} / N \mathbb{Z} \backslash\{\mathbf{0}\}$.

We now consider a random walk on $\mathbb{Z} / N \mathbb{Z}$ which can jump anywhere: $G_{\nu}=\mathbb{Z} / N \mathbb{Z} \backslash\{\mathbf{0}\}$. Here, we focus for simplicity on the case $N=4$, which is the first non trivial example, and disintegrate the jump measure $\nu$ as follows:

$$
\nu=\nu(\mathbf{1}) \delta_{\mathbf{1}}+\nu(\mathbf{2}) \delta_{\mathbf{2}}+\nu(\mathbf{3}) \delta_{\mathbf{3}}, \quad \nu(\mathbf{1}) \nu(\mathbf{2}) \nu(\mathbf{3})>0
$$

It can be proven along the same lines as in the previous examples, that a suitable non negative basis for the lattice $\mathscr{L}$ is given by $\mathscr{B}=\left\{\varphi^{*}, \eta^{*}, \xi^{*}\right\}$ where

$$
\varphi^{*}=\mathbb{1}_{1}+\mathbb{1}_{3}, \quad \eta^{*}=4 \mathbb{1}_{1}, \quad \xi^{*}=2 \mathbb{I}_{1}+\mathbb{1}_{2} .
$$

Hence the associated reciprocal characteristics are:

$$
\Phi_{\varphi^{*}}^{\nu}=\nu(\mathbf{1}) \nu(\mathbf{3}), \Phi_{\nu}^{\eta^{*}}=\nu(\mathbf{1})^{4}, \Phi_{\xi^{*}}^{\nu}=\nu(\mathbf{1})^{2} \nu(\mathbf{2})
$$

We now turn our attention to the integral formula (14). Simple functions $F \in \mathscr{H}_{\xi^{*}}$ are of the form:

$$
F\left(X,\left(\gamma_{1}, \gamma_{2}, \gamma_{3}\right)\right)=f(X) \mathbb{I}_{\left\{g_{1}=g_{2}=\mathbf{1}, g_{3}=\mathbf{2}\right\}} h\left(t_{1}, t_{2}, t_{3}\right) .
$$

$\left(\gamma_{1}, \gamma_{2}, \gamma_{3}\right)$ is in the support of $F$ if two jumps of value $\mathbf{1}$ happen at times $t_{1}, t_{2}$ and one jump of value $\mathbf{2}$ at time $t_{3}$, leading to a global null displacement since $\mathbf{4}=\mathbf{0}$. The formula (14) reads:

$$
\mathbb{Q}\left(f(X) \sum_{\substack{\left(t_{1}, t_{2}, t_{3}\right): t_{1} \neq t_{2}, \Delta X_{t_{1}}=\Delta X_{t_{2}}=1 \\ \Delta X_{t_{3}}=\mathbf{2}}} h\left(t_{1}, t_{2}, t_{3}\right)\right)=\nu(\mathbf{1})^{2} \nu(\mathbf{2}) \mathbb{Q}\left(\int_{[0,1]^{3}} f\left(\Psi_{\gamma_{1}, \gamma_{2}, \gamma_{3}} X\right) h\left(t_{1}, t_{2}, t_{3}\right)\right) d t_{1} d t_{2} d t_{3} .
$$

The distribution of the random vector $\left(n_{1}, n_{2}, n_{3}\right)$ under the 00 -bridge is given by the following identities, valid for any $v: \mathbb{N}^{3} \rightarrow \mathbb{R}$ :

$$
\begin{aligned}
& \int v\left(n_{1}, n_{2}, n_{3}\right) n_{1} n_{3} \chi\left(d n_{1}, d n_{2}, d n_{3}\right)=\nu(\mathbf{1}) \nu(\mathbf{3}) \int\left(v\left(n_{1}+1, n_{2}, n_{3}+1\right) \chi\left(d n_{1}, d n_{2}, d n_{3}\right)\right. \\
& \int v n_{1}\left(n_{1}-1\right)\left(n_{1}-2\right)\left(n_{1}-3\right) \chi\left(d n_{1}, d n_{2}, d n_{3}\right)=\nu(\mathbf{1})^{4} \int v\left(n_{1}+4, n_{2}, n_{3}\right) \chi\left(d n_{1}, d n_{2}, d n_{3}\right) \\
& \int v\left(n_{1}, n_{2}, n_{3}\right) n_{1}\left(n_{1}-1\right) n_{2} \chi\left(d n_{1}, d n_{2}, d n_{3}\right)=\nu(\mathbf{1})^{2} \nu(\mathbf{2}) \int v\left(n_{1}+2, n_{2}+1, n_{3}\right) \chi\left(d n_{1}, d n_{2}, d n_{3}\right) \\
& \chi\left(n_{1} \mathbf{1}+n_{2} \mathbf{2}+n_{3} \mathbf{3}=\mathbf{0}\right)=1
\end{aligned}
$$

In this situation, again $\mathbb{P}_{\nu}$ is the unique random walk of its reciprocal class.

### 3.3. The state space is a product group

Consider the product of two groups, say $G$ and $G^{\prime}$, and two non negative finite measures on them, say $\nu$ and $\nu^{\prime}$, such that in both cases (H1) and (H2) are satisfied. Then, the
product group $G \times G^{\prime}$ equipped with the product measure $\nu \otimes \nu^{\prime}$ fulfills (H1) and (H2) too. The key idea is as follows: if $\mathscr{B}$ and $\mathscr{B}^{\prime}$ are suitable basis of $G$ and $G^{\prime}$ then we can define for all $\eta \in \mathscr{B}$,

$$
\varphi_{\eta}: G_{\nu} \times G_{\nu^{\prime}} \rightarrow \mathbb{N}, \quad \varphi_{\eta}\left(g, g^{\prime}\right)=\eta(g)
$$

and for all $\eta^{\prime} \in \mathscr{B}^{\prime}$,

$$
\varphi_{\eta^{\prime}}: G_{\nu} \times G_{\nu^{\prime}} \rightarrow \mathbb{N}, \quad \varphi_{\eta^{\prime}}\left(g, g^{\prime}\right)=\eta^{\prime}\left(g^{\prime}\right)
$$

The set $\mathscr{B}_{\otimes}=\left\{\varphi_{\eta}\right\}_{\eta \in \mathscr{B}} \cup\left\{\varphi_{\eta^{\prime}}\right\}_{\eta^{\prime} \in \mathscr{B}^{\prime}}$ is an appropriate basis for the lattice of skeletons defined on the product group.

Example 30. Random walk on the $d$-dimensional discrete hypercube $(\mathbb{Z} / 2 \mathbb{Z})^{d}$.
The d-dimensional discrete hypercube is the d-product of the cyclic group with two elements. We denote by $\left(e_{1}, \cdots, e_{d}\right)$ its canonical basis.
A random walk on the hypercube is defined uniquely through its jump measure $\nu=$ $\sum_{i=1}^{d} \nu(i) \delta_{e_{i}}$. Since it can be realized as the product of $d$ random walks on $\mathbb{Z} / 2 \mathbb{Z}$, the basis $\mathscr{B}:=\left\{\varphi_{i}^{*}\right\}_{1 \leq i \leq d}, \varphi_{i}^{*}=2 \mathbb{1}_{e_{i}}$, is a suitable choice.
For the integral characterization it is enough to consider loops of length $\ell=2$. However, we have here d different skeletons to consider. Test functions of the form

$$
F(X, \gamma)=f(X) \mathbb{I}_{\left\{g_{1}=g_{2}=e_{i}\right\}} h\left(t_{1}, t_{2}\right), \quad 1 \leq i \leq d
$$

belong to $\mathscr{H}_{\varphi_{i}^{*}}$. For any $i \in\{1, \cdots, d\}$ fixed, (14) reads as:

$$
\mathbb{Q}\left(f(X) \sum_{\substack{\left(t_{1}, t_{2}\right): t_{1} \neq t_{2}, \Delta X_{t_{1}}=\Delta X_{t_{2}}=e_{i}}} h\left(t_{1}, t_{2}\right)\right)=\nu(i)^{2} \int_{[0,1]^{2}} \mathbb{Q}\left(f\left(\Psi_{\gamma_{1}, \gamma_{2}} X\right)\right) h\left(t_{1}, t_{2}\right) d t_{1} d t_{2} .
$$

Concerning the distribution of the random vector $\left(n_{e_{1}}, \cdots, n_{e_{d}}\right)$, it has independent marginals $\chi_{i}, 1 \leq i \leq d$, which are characterized through the system of equations: for all $v \in \mathcal{B}^{+}(\mathbb{N})$,

$$
\begin{aligned}
\int v(n) n(n-1) \chi_{i}(d n) & =\nu(i)^{2} \int v(n+1) \chi_{i}(d n) \\
\chi_{i}(n \in 2 \mathbb{N}) & =1
\end{aligned}
$$

## Acknowledgments

The authors are very grateful to Paolo Dai Pra and Hans Zessin for fruitful discussions on that topics. The first author acknowledges the Berlin Mathematical School and the Research Training Group 1845 Stochastic Analysis with Applications in Biology, Finance and Physics for their financial support.

## References

[1] S. Bernstein. Sur les liaisons entre les grandeurs aléatoires. In Vehr. des intern. Mathematikerkongress Zürich, volume I. 1932.
[2] J.M.C. Clark. A local characterization of reciprocal diffusions. Applied Stochastic Analysis, 5:45-59, 1991.
[3] G. Conforti, P. Dai Pra, and S. Rœlly. Reciprocal class of jump processes. Preprint, available at http://opus.kobv.de/ubp/volltexte/2014/7077/pdf/premath06.pdf, 2014.
[4] G. Conforti and C. Léonard. Reciprocal classes of random walks on graphs. Preprint, 2014.
[5] G. Conforti, R. Murr, C. Léonard, and S. Rolly. Bridges of Markov counting processes. Reciprocal classes and duality formulae. Electronic Communications in Probability, 2014.
[6] B. Jamison. Reciprocal processes: The stationary Gaussian case. The Annals of Mathematical Statistics, 41(5):1624-1630, 1970.
[7] B. Jamison. Reciprocal processes. Zeitschrift für Wahrscheinlichkeitstheorie und Verwandte Gebiete, 30(1):65-86, 1974.
[8] B. Jamison. The Markov processes of Schrödinger. Zeitschrift für Wahrscheinlichkeitstheorie und Verwandte Gebiete, 32(4):323-331, 1975.
[9] C. Léonard. A survey of the Schrödinger problem and some of its connections with optimal transport. Discrete Contin. Dyn. Syst. A, 34(4):1533-1574, 2014.
[10] C. Léonard, S. Rœlly, and J.-C. Zambrini. Reciprocal processes.a measuretheoretical point of view. Probability Surveys, 11:237-269, 2014.
[11] J. Mecke. Stationäre zufällige Maße auf lokalkompakten Abelschen Gruppen. Zeitschrift für Wahrscheinlichkeitstheorie und Verwandte Gebiete, 9(1):36-58, 1967.
[12] R. Murr. Reciprocal classes of Markov processes. An approach with duality formulae. PhD thesis, Universität Potsdam, 2012. Available at opus.kobv.de/ubp/volltexte/2012/6301/pdf/premath26.pdf.
[13] R. Murr. Characterization of infinite divisibility by duality formulas. application to lévy processes and random measures. Stochastic Processes and their Applications, 123(5):1729-1749, 2013.
[14] N. Privault. Stochastic Analysis In Discrete and Continuous Settings. Springer Berlin Heidelberg, 2009.
[15] S. Roelly. Reciprocal processes. A stochastic analysis approach. In V. Korolyuk, N. Limnios, Y. Mishura, L. Sakhno, and G. Shevchenko, editors, Modern Stochastics and Applications, volume 90 of Optimization and Its Applications, pages 53-67. Springer, 2014.
[16] S. Roelly and M. Thieullen. Duality formula for the bridges of a brownian diffusion: Application to gradient drifts. Stochastic Processes and their Applications, 115(10):1677-1700, 2005.
[17] L. Saloff-Coste. Random walks on finite groups. Probability on discrete structures, Encyclopaedia Math. Sci.,110:263-346. Springer,Berlin, 2004.
[18] E. Schrödinger. Über die Umkehrung der Naturgesetze. Sitzungsberichte Preuss.

Akad. Wiss. Berlin. Phys. Math., 144:144-153, 1931.
[19] M. Thieullen. Second order stochastic differential equations and non-Gaussian reciprocal diffusions. Probability Theory and Related Fields, 97(1-2):231-257, 1993.
[20] M. Thieullen and J.-C. Zambrini. Symmetries in the stochastic calculus of variations. Probability Theory and Related Fields, 107(3):401-427, 1997.

